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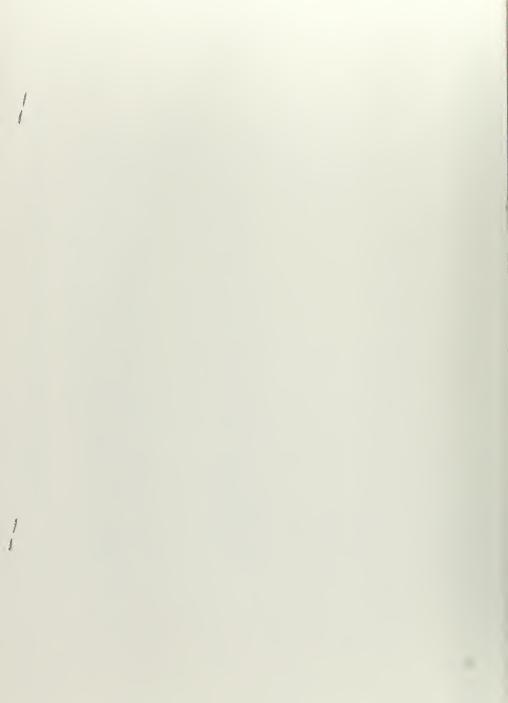
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AN ANALYSIS OF SINGLE STAGE AXIAL-FLOW TURBINE PERFORMANCE USING THREE-DIMENSIONAL

CALCULATING METHODS

ROBERT GLEN HARRISON







#### AN ANALYSIS OF SINGLE STAGE

AXIAL-FLOW TURBINE PERFORMANCE USING

THREE-DIMENSIONAL CALCULATING METHODS

by

Robert Glen Harrison Lieutenant, United States Navy B.M.E., University of Louisville, 1959

Submitted in partial fulfillment of the requirements for the degree of

AERONAUTICAL ENGINEER

from the

NAVAL POSTGRADUATE SCHOOL September 1967 I E PER P.

ABSTRACT

The method of turbine performance prediction developed by Vavra and Eckert has been refined in this analysis to realize more of the potential of the three-dimensional calculating methods. Mach number and rotor tip clearance effects on blade outlet angles and loss coefficients have been localized rather than averaged over the blade height. An approximation for streamline curvature has been used.

Performance curves were determined for two single stage axialflow turbines located at the Propulsion Laboratory of the Naval Postgraduate School. Test results were available for one of the turbines. Agreement between predicted and experimental performance values was generally within 3 per cent. Thesis by Robert G. Harrison entitled: "An Analysis of Single-Stage Axial-Flow Turbine Performance Using Three-Dimensional Calculating Methods"

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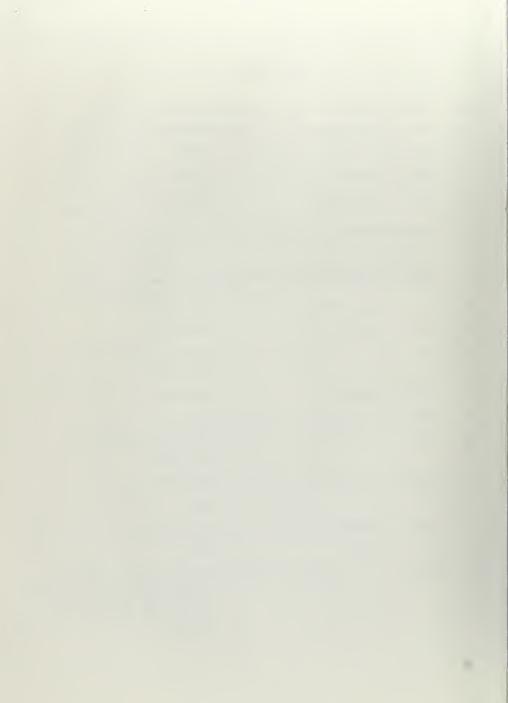
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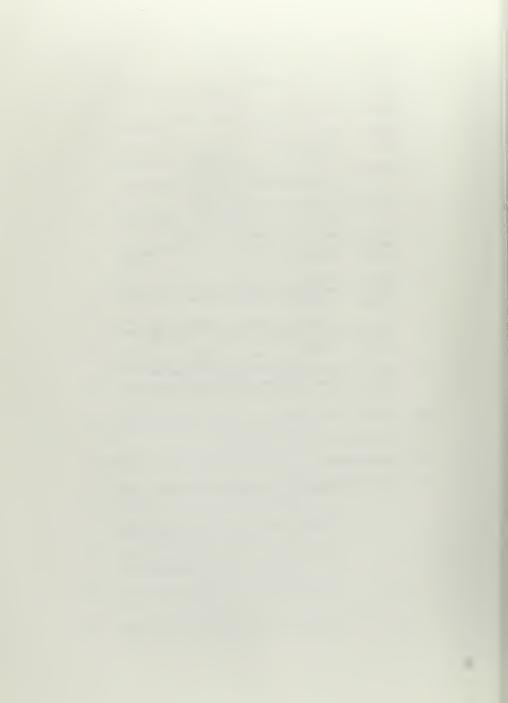
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# TABLE OF SYMBOLS

Symbols	
A	Area (in. <sup>2</sup> )
a	Throat opening of blade channel (in.)
В	Shroud factor used in rotor loss coefficient calculations (dimensionless)
b	Blade's departure from being straight-backed (in.)
c <sub>1</sub>	Conversion factor, $2gJ(ft^2-1b_m/sec^2-BTU)$
С	Blade Chord (in.)
c <sub>p</sub>	Specific heat, constant pressure (BTU/ $1b_m^{-0}R$ )
É	Kinetic energy rate (ft-lb/sec)
e	Mean radius of curvature of back of blade (in.)
g	Universal gravitational constant (32.174 $1b_m$ -ft/1b-sec <sup>2</sup> )
Н	Total enthalpy (BTU/1b <sub>m</sub> )
H***	Boundary layer energy parameter ( 5 * )
h	Static enthalpy (BTU/1b <sub>m</sub> )
h	Blade height (in.)
HP	Horsepower
I	Integrand
i	Incidence angle (deg. or radians)
ī	Unit vector
is	Stalling incidence angle (deg. or radians)
J	Conversion factor (778.16 ft-1b/ETU)
j	Distance from throat to trailing edge of blade (in.)
k	Tip clearance (in.)
k	Isentropic head coefficient (dimensionless)

```
Symbols
  L
            Distance between stations 0 and 1 and between
            stations 1 and 2 (in.)
  M
            Mach number
  M
            Moment (ft-1b)
            Mass flowrate (slugs/sec)
  m
  m
            Exponent used in boundary layer calculations,
            see Eq. 146 (dimensionless)
  P
            Pressure (psia)
  R
            Gas constant for air (53.345 ft-1b/1b_-OR)
            Radius (in.)
  r
  r*
            Theoretical degree of reaction (dimensionless)
            Blade spacing (in.)
            Entropy (BTU/1b - R)
  8*
            Non-dimensional entropy (
            Temperature (OR)
  т
            Blade thickness (in.)
            Blade trailing edge thickness (in.)
  t
  t*
            Projection of blade trailing edge thickness on the
            exit plane of the blade row (in.)
  U
            Peripheral velocity (ft/sec)
  11
            Velocity within a boundary layer (ft/sec)
  ٧
            Absolute velocity (ft/sec)
  W
            Relative velocity (ft/sec)
            Weight flowrate (lbm/sec)
  Wf
            Fraction of the total flowrate which passes between
            the hub and any other streamline (dimensionless)
            Reference flowrate (in. 2)
  Wref
```

### Symbols

- X Non-dimensional radius  $\underline{r}$  where  $r_m$  is the mean streamline radius  $\underline{r}_m$
- X Shroud factor for calculation of rotor outlet angles (dimensionless)
- X See Eq. 145
- Y Non-dimensional axial velocity  $\frac{V_A}{V_{A_m}}$  where  $^V_{A_m}$  is the mean streamline axial velocity
- Y Pressure loss parameter, see Eq. 167
- y Distance from wall of a point in a boundary layer (in.)
- Z Number of blades

#### Greek Letters

- Absolute gas flow angles (deg. or radians)
- B Gas flow angles relative to rotor (deg. or radians)
- $\mathcal{B}_{a}$  Blade inlet angle (deg. or radians)
- Specific heat ratio (dimensionless)
- $\Delta R$  Streamline displacement, see Fig. 2 (in.)
  - $\delta$  Boundary layer thickness (in.)
  - 8 Referred pressure  $\frac{P_{to}}{14.7}$  (dimensionless)
  - Sr Streamline displacement, see Fig. 2 (in.)
  - S\* Boundary layer displacement thickness (in.)
- $\delta$  \*\*\* Boundary layer energy thickness (in.)
- £ Loss coefficient (dimensionless)
- The Efficiency (dimensionless)
- $\gamma$  Non-dimensional distance from the wall in a boundary layer ( )
- Referred temperature 518.4 (dimensionless)
- K Streamline curvature factor (dimensionless)

#### Greek Letters

Angle between flow and axis of turbine in a meridional plane (deg. or radians)

Factor used in predicting secondary loss coefficients, see Eq. 171 (dimensionless)

Area restriction factor (dimensionless)

 $\rho$  Density  $(1b_m/ft^3)$ 

Non-dimensional flow function

ω Angular velocity (radians/sec)

### Subscripts

A Axial

d Discharge

E Equivalent

H Hub

is Isentropic expansion from total inlet conditions

m Mean streamline

o Station ahead of stator

P Profile

R Relative

r Radial

ref Referred

rea. Required

S Stator

S Shroud

s Isentropic expansion from equivalent total conditions

T Tip

t Total

# Subscripts

- th Theoretical
- u Tangential
- a Axial direction, cyclindrical coordinates
- $\theta$  Peripheral direction, cyclindrical coordinates
- 1 Station between stator and rotor
- 2 Station after rotor

# Superscript

\*\* Fefers to predicted values for the mean streamline



### 1. Introduction

Turbines form an important part of propulsion systems. To optimize a design it is necessary to know the performance at off-design conditions as well as the performance at the design point. Since the testing of a prototype is very costly and time consuming, it is of great advantage to be able to predict the performance characteristics by means of theoretical methods. The more advanced a system is the more important it becomes to improve the accuracy of these methods. Hence it is necessary to base these methods on the fundamental laws of fluid dynamics rather than on rule-of-thumb approximations. The latter basis is possible only when data on previous designs are available. For new and advanced configurations, it will be necessary to apply refined methods which enable the designer and systems engineer to predict the effect of proposed design changes.

The principal equations that describe the flow properties in turbomachines are well known. These same equations are used as a basis for all proper "three-dimensional" calculating methods that have been cited in the technical literature. The methods differ however in the manner in which these equations are manipulated and applied.

This thesis is concerned with the refinement of the three-dimensional method of analysis developed by Vavra and Eckert. Based on the physical dimensions of particular test turbines that are available at the Turbo-Propulsion Laboratory of the Naval Postgraduate School, performance curves were determined for these machines. Concurrently with the performance analysis, experimental tests were conducted on one of these turbines so that actual experimental results could be used to judge the accuracy of the proposed theoretical performance evaluation.

In addition to his publications and classroom lectures which constituted the foundation for this analysis, Professor Vavra was very generous in providing guidance and counsel during the period of this work. For this I am greatly appreciative. I would also like to thank Lieutenants P. M. Commons and J. A. Messegee for making their experimental results available.

### 2. Basis for the Analysis

The two conservation equations that must be satisfied to obtain a solution for the flow in a turbine are the equations of motion and continuity. These equations are satisfied at stations between blade rows. The method used is that given by Vavra. Vavra developed the equations of motion and continuity for absolute flows. The equations in this form are readily useable for the position after the stator. Eckert later developed these equations for relative flows to be used for rotor calculations. Eckert's conversion allows relative flow quantities to be handled directly without conversion to an absolute system. Eckert's approach also avoids iteration procedures to determine the total enthalpy after the rotor.

The assumptions made for the development of the equations used in the performance analysis are:

- An infinite number of blades in each row so that downstream effects are not felt upstream.
- Axisymmetric flow at the stations where the equations of motion are solved.
- 3. Adiabatic and steady flow so that the total enthalpy along any given streamline is constant through the stator and the relative total enthalpy is constant through the rotor.
- 4. All entropy changes are assumed to occur in the blade channels that are located ahead of the stations where the equations of motion are satisfied. Hence at the calculating stations the flow is assumed to be isentropic along particular streamlines.

With the above assumptions, the equations of motion for absolute and relative flows, respectively, are

$$\nabla H = \nabla \times (\nabla \times \nabla) + T \nabla S \tag{1}$$

$$\nabla H_{R} = \overline{W} \times (\nabla \times \overline{W} + 2 \overline{\omega}) + T \nabla S$$
 (2)

<sup>&</sup>lt;sup>1</sup>Vavra, M. H., <u>Aero-Thermodynamics and Flow in Turbomachines</u>. New York, London: John Wiley and Sons, Inc.; 1960, Chapter 16.

<sup>&</sup>lt;sup>2</sup>Eckert, R. H., Performance Analysis and Initial Tests of a Transonic Turbine Test Rig (USNPGS Thesis, May 1966), pp. 149-155.

Relative total enthalpy can be written as

$$H_R = h_1 + \frac{W_1^2}{2} - \frac{U_1^2}{2} = h_2 + \frac{W_2^2}{2} - \frac{U_2^2}{2} = h_{25} + \frac{W_{2th}^2}{2} - \frac{U_2^2}{2}$$
 (3)

Equivalent enthalpy is defined as

$$H_{E} = h_{2S} + \frac{W_{z+h}^{2}}{2} = H_{R} + \frac{U_{z}^{2}}{2}$$
 (4)

Similar to  $H_{\rm R}$ , the equivalent enthalpy  $H_{\rm E}$  is constant along a streamline for the adopted assumptions. Equivalent enthalpy can also be written as

e written as
$$H_{E} = h_{2} + \frac{W_{z}^{2}}{2} = h_{1} + \frac{W_{1}^{2}}{2} + \frac{U_{2}^{2} - U_{1}^{2}}{2}$$
(5)

The introduction of the equivalent enthalpy allows this quantity to be used for the rotor in a manner analogous to the way total enthalpy is used for the stator.

For equations used in this analysis, the subscripts refer to:

0 - station ahead of stator

1 - station ahead of rotor

2 - station after rotor

is - isentropic expansion from Pt

s - isentropic expansion from  $P_{t_n}$ 

th - the theoretical value

Figure 1 is a temperature-entropy diagram showing the thermodynamic process along a particular streamline for a single stage turbine. In general, the fluid properties will vary from streamline to streamline. The method in which the loss coefficients are applied is also indicated in Fig. 1. The loss coefficients are defined as

$$J_{s} = \frac{V_{l_{th}}^{2} - V_{l}^{2}}{V_{l_{th}}^{2}} \qquad \text{For the Stator}$$
 (6)

$$U_{R} = \frac{W_{eth}^2 - W_{eth}^2}{W_{eth}^2}$$
 For the Rotor (7)

The coordinate system that will be used in the analysis is indicated in Fig. 2. This figure also shows the general layout of the type turbine to which the prediction performance analysis is applied.

Fign convention for the various angles that are needed in the analysis 1 indicated in Fig. 3.

The modification of Eq. 2 into a form that can be used for the analysis is given in Appendix A, Section 1. The appropriate form of Eq. 1 can be obtained from the modification of Eq. 2 if the angular velocity  $\omega$  is set equal to zero. Other differences of the final equations derived from Eqs. 1 and 2 are listed in Appendix A,

Section 1. Equation 1 can then be written
$$\frac{d(\mathcal{L}_{N}Y_{i}^{2})}{dX_{i}} = -\cos^{2}\alpha_{i}\left[\left(-K2Y_{m}\frac{SY}{L^{2}}\right) - \left(\frac{L^{2}+\left(NR/2\right)^{2}}{L^{2}}\right)\frac{dS_{i}^{*}}{dX_{i}}\right] - 2 TAN\alpha_{i}\frac{d\alpha_{i}}{dX_{i}}$$

$$-\frac{2}{X_{i}}S_{i}N^{2}\alpha_{i} + \frac{C_{i}C_{i}C_{i}^{2}\alpha_{i}}{Y_{i}^{2}V_{A_{i}m}^{2}}\frac{dH}{dX_{i}} - \left[\frac{C_{i}H\cos^{2}\alpha_{i}}{Y_{i}^{2}V_{A_{i}m}^{2}} - S_{i}N^{2}\alpha_{i}\right]\frac{dS_{i}^{*}}{dX_{i}}$$
(8)

Equation 2 becomes  $\frac{d(\sqrt{M} \frac{\gamma_2^2}{2})}{d X_2} = -\cos^2 \beta_2 \left[ \left( K2 r_m \frac{\delta r}{L^2} \right) - \left( \frac{L^2 + \left( \frac{\delta R}{2} \right)^2}{L^2} \right) \frac{d S_2^*}{d X_2} \right] - 2 TAN \beta_2 \frac{d \beta_2}{d X_2} - \frac{2}{X_2} SIN^2 \beta_2$ 

$$-\frac{4 U_{m} \cos k_{1} \sin k_{2}}{W_{A_{m}} Y_{2}^{2}} + \frac{C_{1} \cos^{2} k_{2}}{W_{A_{m}}^{2} Y_{2}^{2}} \frac{d H_{E}}{d X_{2}} - \left[ \frac{C_{1} H_{E} \cos^{2} k_{2}^{2} - \sin^{2} k_{2}^{2}}{W_{A_{m}}^{2} Y_{2}^{2}} \frac{d S_{2}^{*}}{d X_{2}} \right] \frac{d S_{2}^{*}}{d X_{2}}$$
(9)

where:

$$Y = \frac{V_A}{V_{A_{m}}} \text{ or } \frac{W_A}{W_{A_{m}}}$$
 for Eqs. 8 and 9 respectively (subscript m refers to mean streamline) 
$$X = \frac{Y_A}{Y_{A_{m}}}$$
 
$$X = \frac{Y_A}{Y_{A_{m}}}$$
 streamline curvature factor streamline displacements shown in Fig. 2 
$$\Delta R$$
 
$$L = \frac{L_1 + L_2}{2}$$
 (see Fig. 2) 
$$S^* = \frac{S}{C_P}$$
 
$$C_1 = 2.9 \text{ J}$$

The equation of continuity is used in its non-dimensional form by introducing a flow function  $\Phi$ . Development of  $\Phi$  is given in Appendix A, Section 2. In differential form the equations for absolute and relative flows can then be expressed by

$$\frac{d\dot{w}\sqrt{1_{to}}\sqrt{\frac{R}{9}}}{P_{to}} = dA = dA\sqrt{\frac{28}{1_{to}}\left[\left(\frac{P}{P_{to}}\right)^{\frac{2}{8}} - \left(\frac{P}{P_{to}}\right)^{\frac{4+1}{8}}\right]}$$
(10)

and

$$\frac{d\hat{W}_{t_{\epsilon}}}{P_{t_{\epsilon}}} \sqrt{\frac{R}{9}} = dA \vec{q} = dA \sqrt{\frac{28}{8-1} \left[ \left( \frac{P}{P_{t_{\epsilon}}} \right)^{\frac{2}{8}} - \left( \frac{P}{P_{t_{\epsilon}}} \right)^{\frac{4+1}{8}} \right]}$$
(11)

The differential element of area dA is

$$JA = \xi \, \Xi \, a \, dr \tag{12}$$

where:

2 = number of blades

a = blade exit opening (see Fig. 46)

= are: restriction coefficient

Since  $\bar{\Phi}$  is valid for isentropic flow only, the restriction factor  $\bar{\xi}$  must be introduced to correct the actual flow area to an effective area which accounts for the restrictions due to the boundary layers on both sides of the flow channel.

The factor 
$$\xi$$
 can be expressed by  $\frac{\xi}{\xi} = 1 - \xi \frac{\xi^*}{q} = \frac{11^{***} - 1}{11^{***} - 1 + \frac{11}{2}}$  (13)

In Eq. 13, 5 \* is the boundary layer displacement thickness and h\*\*\* is the so-called energy parameter defined as the energy thickness divided by the displacement thickness. The term - represents the loss that is assumed to occur from the inlet to the throat of the blade channel, where y is the loss coefficient representing all the losses across the row of blades. The profile loss coefficient was used by Eckert 4 to represent the loss prior to the blade throat. Percentage of the total loss due to profile losses will vary considerably depending on blade geometry, radial position, and the incidence of the flow on the leading edge of the blade. Since secondary flow and tip clearance effects result in losses in the blade channel, half the total loss coefficient provides a better average representation of the losses in the blade channel prior to the throat. The basis for the development of and H\*\*\* as used is given in Appendix A, Sections 3 and 4.

By multiplying and dividing by  $a_m \cdot a_m = Eq. 12$  is

$$1/r = \frac{1}{2} - \frac{1}{2} A^{X} \cdot a_{m} t_{m}$$
 (14)

Vavra, M. H., Problems of Fluid Mechanics in Radial Turbomachines (Rhode-Saint-Genese, Belgium: Von Karmán Institute for Fluid Dynamics, 1965) VKI Course Note 55b, pp. G46-50.

ackert, op. cit., p. 44.

After integration Eqs. 10 and 11 become, respectively,

$$\frac{\overrightarrow{W}\sqrt{T_{ro}}}{P_{ro}}\sqrt{\frac{R}{g}} = a_m Z r_m \int_{a_m}^{A_r} \xi \int dX$$
 (15)

and

$$\frac{\text{WVT}_{te}}{f_{te}^{2}}\sqrt{\frac{2}{9}} = a_{m} Z r_{m} \int_{X_{t}}^{X_{T}} \frac{a}{a_{m}} \xi J dX$$
 (16)

The flowrate  $\dot{w}$  can be computed from the conditions ahead of the stator. Then a reference flowrate is defined by

$$W_{ref} = \left[ \frac{\dot{w} \sqrt{T_{to}} \sqrt{R}}{P_{to}} \sqrt{\frac{R}{g}} \right]$$
 (17)

where  $\boldsymbol{W}_{\mbox{ref}}$  is in square inches. Continuity will be satisfied for the stator by

$$\left[ \mathcal{A}_{m} \stackrel{?}{=} r_{m} \int_{\chi_{H}}^{\chi_{T}} \xi \, \underbrace{1}_{d} \, d\chi \right]_{STATCR} = W_{ref}$$
(18)

Similarly for the rotor

$$\left| a_{m} \neq r_{m} \right\rangle \frac{X_{T}}{R_{to}} \sqrt{\frac{T_{to}}{T_{tr}}} \frac{a}{a_{m}} \xi \int dX = W_{rcf}$$
(19)

The in luence of the leakage flow through the radial tip clearence has not been accounted for in Eq. 19.

The element of area between the blade tips and the shroud is

$$dA = 2\pi r dr$$
 (20)

The flow through the tip clearance area is

$$\left[\frac{\sqrt{1-z}}{P_{FE}}\right]_{T,P} = 2\pi a_m Z r_m^2 \int_{X_T}^{X_S} \frac{X}{Z a_m} \xi \int_{X_T}^{X_S} dX$$
(21)

Since the tip clearance is relatively small, the values of

 $\frac{\pi}{4}$ ,  $\frac{\pi}{4}$ ,  $\frac{\pi}{4}$ ,  $\frac{\pi}{4}$ , and  $\frac{\pi}{4}$  for the tip will be used in Eq. 21.

With these assumptions Eq. 19 can be expressed by 
$$\left[ \int_{\chi_H}^{\chi_T} \frac{\chi_T}{P_{fo}} \sqrt{\frac{T_{fo}}{T_{fc}}} \frac{\alpha}{\alpha_m} \xi \, \frac{1}{2} \, dX + 2\pi r_m \int_{\chi_T}^{\chi_S} \frac{F_{fo}}{P_{fo}} \sqrt{\frac{T_{fo}}{T_{fc}}} \, \xi \, \frac{1}{2} \, \chi \right)_{\tau_1 F} \frac{d\chi}{z_{\alpha_m}} \right]_{R \in To R}$$

 $= \frac{\text{Wre4}}{\left[\alpha_{m} \pm \gamma_{m}\right]_{\text{RoToR}}} \tag{32}$  The assumptions used to arrive at Eq. 22 are obviously incorrect in

The assumptions used to arrive at Eq. 22 are obviously incorrect in two respects. First, the flow represented by  $\frac{1}{2}$  is not

perpendicular to the tip clearance area. Second, the effective area represented by  $\xi_{\rm T}$  is larger than that which probably occurs because of the relatively large boundary layers that exist on the shroud and blade tips. The exact behavior of the flow in the small region between the shroud and the blade tips is impossible to predict without further tests. However, it is felt that the flowrate through this space as represented in Eq. 22 is too large for the reasons just mentioned. Therefore a more accurate approximation of this flowrate will be obtained if the last term on the left side of Eq. 22 is

divided by 2, yielding
$$\begin{bmatrix}
\sqrt{\frac{x_{r_{e}}}{R_{to}}} \sqrt{\frac{T_{to}}{T_{tr}}} & \frac{a}{a_{rx}} \xi & dX + \pi v_{rx} \sqrt{\frac{x_{r_{e}}}{R_{to}}} \sqrt{\frac{T_{to}}{T_{te}}} \xi & dX
\end{bmatrix}$$
ROTOR

$$= \frac{\text{Wref.}}{\left[\alpha_{m} Z r_{rn}\right]_{ROTOR}}$$
 (23)

The tip clearance flow included in Eq. 23 can be obtained by integration

$$\pi r_{m} \int_{X_{-}}^{X_{-}} \left\langle \frac{R_{c}}{R_{o}} \sqrt{\frac{T_{to}}{T_{tc}}} \xi \right| dX \right)_{T,P} \frac{dX}{Za_{m}} = \frac{\pi r_{m}}{Za_{m}} \left( \frac{R_{c}}{R_{o}} \sqrt{\frac{T_{to}}{T_{tc}}} \xi \right) X \right)_{T,P} \left[ X_{s} - X_{T} \right]$$

$$= \frac{\pi r_{m}}{Za_{m}} \left( \frac{R_{c}}{R_{o}} \sqrt{\frac{T_{to}}{T_{tc}}} \xi \right) X \right)_{T,P} \frac{k}{r_{m}} = \frac{\pi k r_{T}}{a_{m} Z r_{m}} \left( \frac{R_{c}}{R_{c}} \sqrt{\frac{T_{to}}{T_{tc}}} \xi \right) X \right)_{T,P} \left[ X_{s} - X_{T} \right]$$

$$= \frac{\pi r_{m}}{Za_{m}} \left( \frac{R_{c}}{R_{o}} \sqrt{\frac{T_{to}}{T_{tc}}} \xi \right) X \right)_{T,P} \left[ X_{s} - X_{T} \right]$$

$$= \frac{\pi r_{m}}{Za_{m}} \left( \frac{R_{c}}{R_{c}} \sqrt{\frac{T_{to}}{T_{tc}}} \xi \right) X \right)_{T,P} \left[ X_{s} - X_{T} \right]$$

$$= \frac{\pi r_{m}}{Za_{m}} \left( \frac{R_{c}}{R_{c}} \sqrt{\frac{T_{to}}{T_{tc}}} \xi \right) X \right)_{T,P} \left[ X_{s} - X_{T} \right]$$

$$= \frac{\pi r_{m}}{Za_{m}} \left( \frac{R_{c}}{R_{c}} \sqrt{\frac{T_{to}}{T_{tc}}} \xi \right) X \right)_{T,P} \left[ X_{s} - X_{T} \right]$$

$$= \frac{\pi r_{m}}{Za_{m}} \left( \frac{R_{c}}{R_{c}} \sqrt{\frac{T_{to}}{T_{tc}}} \xi \right) X \right)_{T,P} \left[ X_{s} - X_{T} \right]$$

$$= \frac{\pi r_{m}}{Za_{m}} \left( \frac{R_{c}}{R_{c}} \sqrt{\frac{T_{to}}{T_{tc}}} \xi \right) X \right)_{T,P} \left[ X_{s} - X_{T} \right]$$

$$= \frac{\pi r_{m}}{Za_{m}} \left( \frac{R_{c}}{R_{c}} \sqrt{\frac{T_{to}}{T_{tc}}} \xi \right) X \right)_{T,P} \left[ X_{s} - X_{T} \right]$$

$$= \frac{\pi r_{m}}{Za_{m}} \left( \frac{R_{c}}{R_{c}} \sqrt{\frac{T_{to}}{T_{tc}}} \xi \right) X \right)_{T,P} \left[ X_{s} - X_{T} \right]$$

$$= \frac{\pi r_{m}}{Za_{m}} \left( \frac{R_{c}}{R_{c}} \sqrt{\frac{T_{to}}{T_{tc}}} \xi \right) X \right)_{T,P} \left[ X_{s} - X_{T} \right]$$

$$= \frac{\pi r_{m}}{Za_{m}} \left( \frac{R_{c}}{R_{c}} \sqrt{\frac{T_{to}}{T_{tc}}} \xi \right) X \right)_{T,P} \left[ X_{s} - X_{T} \right]$$

$$= \frac{\pi r_{m}}{Za_{m}} \left( \frac{R_{c}}{R_{c}} \sqrt{\frac{T_{to}}{T_{tc}}} \xi \right) X \right)_{T,P} \left[ X_{s} - X_{T} \right]$$

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$$= \frac{\pi r_{m}}{Za_{m}} \left( \frac{R_{c}}{R_{c}} \sqrt{\frac{R_{c}}{T_{tc}}} \xi \right) X \right)_{T,P} \left[ X_{s} - X_{T} \right]$$

$$= \frac{\pi r_{m}}{Za_{m}} \left( \frac{R_{c}}{R_{c}} \sqrt{\frac{R_{c}}{T_{tc}}} \xi \right) X \right)_{T,P} \left[ X_{s} - X_{T} \right]$$

$$= \frac{\pi r_{m}}{Za_{m}} \left( \frac{R_{c}}{R_{c}} \sqrt{\frac{R_{c}}{T_{c}}} \frac{R_{c}}{R_{c}} \sqrt{\frac{R_{c}}{T_{c}}} \frac{R_{c}}{T_{c}} \frac{R_{c}}{T_{c}} \frac{R_{c}}{T_{c}} \frac{R_{c}}{T_{c}} \frac{R_{c}}{T_{c}} \frac{R_{c}}{T_{c}} \frac{R_{c}}{T_{c}} \frac{R_{c}}{T_{c}} \frac{R_{c}}{T_{$$

3. Technique for Obtaining Solution

With the equations of motion and continuity in the forms given by Eqs. 8, 9, 18 and 25, a method has been developed to analyze single stage axial turbines, in particular, those available for test in the Turbine Test Rig of the Turbo-Propulsion Labora ory of the Naval Postgraduate School. The method of analysis predicts turbine performance for specified values of inlet total pressure, inlet total temperature, rotor speed and the ratio of total inlet to static discharge pressure  $P_t/P_2$ . This method is similar to that described by Eckert. 5 However, Eckert's analysis neglected some effects which

<sup>5</sup> Ibid, Section 3.

have been accounted for in this developement. Some significant changes made in this method are listed below:

- Stator and rotor outlet angles for a particular radial location are computed using a calculated Mach number for that location rather than an assumed Mach number or the Mach number of the mean streamline.
- Variation of loss coefficients due to changes in blade geometry in the radial direction is accounted for.
- 3. The influence of rotor tip clearance on the rotor outlet angles and loss coefficients is concentrated near the tip of the blade rather than averaging these effects over the full blade height.
- 4. Fourth order polynomials are used to better approximate the curves respresenting blade characteristics as a function of radius and the curves of rotor loss coefficients as a function of incidence.
- Streamline curvature effects have been accounted for in the solution of the equations of motion.

In addition to the assumptions that were mentioned in Section 2 for the development of the particular form of the equation of motion, the conditions ahead of the stator are assumed to be uniform; that is, the total temperature, velocity, and entropy are assumed constant and the flow axial in direction. It is realized that completely uniform conditions are difficult to obtain, but any other assumpt on would be extremely difficult to develop mathematically.

Direct solution of the equations of motion is not possible since they are nonlinear in the dependent variable Y. Likewise, no direct method is possible to satisfy continuity. Solutions of these equations must therefore be gained by making initial assumptions for the values of the axial velocities which must be improved by successive iterations until the equations are satisfied. To account for streamline curvature and slope, a complete solution of the flow through the stator and rotor must first be made by neglecting the effects of curvature in order to determine streamline locations. Then the iteration to account for these effects may progress. These requirements make the use of a high-speed computer a necessity.

This analysis has been programmed for the IBM 360 computer using FORTRAN IV. The program is described in Appendix C. The following paragraphs set forth the procedural steps of the program. The equations are listed in general form without referring to specific streamline locations. In the interest of clarity, however, some relationships will be written in forms similar to those used in the program. For example,  $P_i(2)/P_{to} = \left[ T_{iis}(2)/T_{to} \right]^{\frac{\delta}{\delta-1}}$  will represent the isentropic relationship for the number 2 streamline.

Five streamlines are utilized for the analysis with the number 1 streamline located at the hub and the number 5 streamline located at the tip as shown in Fig. 2. The number 3 streamline will be used as the mean streamline, and the radius of this streamline will be referred to as the mean radius. The radial locations of the streamlines ahead of the stator will be such that the mass flowrate between adjacent streamlines is 25 per cent of the total flowrate. Positions for the streamlines after the stator and after the rotor are initially assumed. The locations of streamlines 2, 3, and 4 then vary during the solution as necessary so that the percentage of the total flowrate between adjacent streamlines does not change. This continuity requirement will be called streamline continuity.

Besides the radii, sufficient input information must be used to effectively reflect the physical characteristics of the stator and rotor blading. Some of the physical properties are introduced directly; such as, the number of stator blades, the number of rotor blades, and the rotor tip clearance. The other quantities used which reflect blade characteristics are throat opening dimensions for the blade channels, discharge angles, rotor blade inlet angles, loss coefficients, and stalling incidences for the rotor.

Throat opening dimension "a" is a function of radius. The best method for introducing this characteristic into the analysis is to enter the measured values of "a" together with the corresponding radii. Then, utilizing the method of least squares, a fourth order polynominal curve is fitted through these points. From the resulting polynomial, the value of "a" for any radius required by streamline continuity can be determined.

Discharge angles are predicted by using a combination of the methods of Vavra and Ainley. Outlet angles are first calculated using the formula which Vavra established from the experimental data of Beer. These angles are then corrected for tip clearance, blade cu.vature, and Mach number effects with the methods given by Ainley. Stator and rotor discharge angles are predicted at three radii; namely, the hub, mean radius, and tip. The method used for calculating these angles is explained in Appendix B, Section 1.

Values of stator gas outlet angles  $\alpha_1$  for the mean streamline are determined for Mach numbers M<sub>1</sub> of 0.5, 0.7, 0.75, 0.8, and 1.0. These values are represented by two parabolic curves of the form

$$\alpha_1 = \alpha + b M_1 + c M_1^2$$
 (26)

The first curve is used for Mach numbers  $M_1$  from 0.5 to 0.75 and the second established interim values of  $\alpha_1$  for  $M_1$  between 0.75 and 1.0. From these curves the flow angle  $\alpha_1$  for the mean streamline can be found for any Mach number  $M_1$ .

The flow angle  $\alpha_1$  is also a function of radius  $r_1$ . Therefore the changes of  $\alpha_1$  for the hub and tip with reference to the mean radius, called  $\Delta\alpha_H$  and  $\Delta\alpha_T$  respectively, must be used. The flow angle  $\alpha_1$  can then be determined for any  $r_1$  by assuming a linear variation or  $\alpha_1$  between the hub and the mean radius and between the mean and the tip. With this assumption and using the Mach number  $M_1$  in Eq. 26 corresponding to the number 2 streamline, the flow angle  $\alpha_1$  for this streamline would be

$$\alpha_{1}(z) = \alpha_{1}^{*} + \frac{\Gamma(z) - \Gamma_{1}^{*}}{\Gamma(1) - \Gamma_{1}^{*}} \Delta \alpha_{H}$$
 (27)

Ainley, D. G. and Mathieson, G. C. R., A Method of Performance Estimation for Axial-Flow Turbines. Aeronautical Research Council, I. & M. No. 2974, 1957. pp. 3-4.

<sup>&</sup>lt;sup>7</sup>Beer, R., Aerodynamic Design and Estimated Performance of a Two-Stage Curtis Turbine for the Liquid Oxygen Turbopump of the M-1 Engine. NASA CR 54764 (AGC 8800-12), Nov. 19, 1965. p. 29.

The superscript \*\* is used with  $r_1$  and  $\mathcal{A}_1$ , in Eq. 27 to indicate the radius initially assumed for the mean streamline and the computed for that radius. Equation 27 can then be used throughout the analysis even though the radial location of the mean streamline may change due to streamline continuity requirements. A similar approach is used to establish the flow angle  $\beta_2$  at the rotor discharge.

The rotor blade inlet angles  $\beta_0$  are measured from the manufacturing drawings of the blade profiles. Using the values of  $\beta_0$  for the hub, mean, and tip streamlines, a parabolic curve is determined which gives  $\beta_0$  as a function of  $X_1$ .

Loss coefficients and stalling incidences are predicted by using the methods of Ainley. <sup>8</sup> For the present method, the stalling incidence is is defined as that at which the loss coefficient is twice the value of the minimum loss coefficient. Following Ainley's methods, loss coefficients are computed as a function of the ratio of flow incidence to stalling incidence  $\frac{\zeta}{\zeta S}$ . Loss coefficients are also a function of blade geometry or radius. Since the stator has zero incidence, its loss coefficient is a function of the radius only.

Stator and rotor loss coefficients and rotor stalling incidences are calculated for the hub, mean, and tip radial locations. Rotor loss coefficients  $\mathcal G$  are determined for values of  $\frac{\dot{\mathcal L}}{\dot{\mathcal L}S}$  ranging from -2.0 to 1.6. Curves of  $\mathcal G_{\mathcal R}$  vs.  $\frac{\dot{\mathcal L}}{\dot{\mathcal L}S}$  are drawn for each of the three radial locations, and values of  $\mathcal G_{\mathcal R}$ , along with the corresponding quantities  $\frac{\dot{\mathcal L}}{\dot{\mathcal L}S}$ , are used to determine fourth order polynomials which approximate these curves. A similar procedur is followed to obtain a fourth order polynomial representing stalling incidence  $i_g$  as a function of radius  $\mathcal T_{\mathcal L}$ . Sample calculations for the prediction of stator and rotor loss coefficients and stalling incidences are contained in Appendix B, Section 2.

<sup>8</sup> Ainley, op. cit., pp. 4-5.

Variation of loss coefficients with radial location is accounted for by assuming a linear variation of these quantities between the hub and the mean radius and between the mean and the tip. To demonstrate the procedure followed in computing rotor loss coefficients, the following example is given. For a particular incidence, the first step in the determination of  $\mathcal{L}_{R}$  for the number 2 streamline would be to calculate  $i_{g}$  by using the radius  $r_{i}$  of that streamline and the polynomial of the form  $i_{g} = i_{g}(r_{1})$ . Loss coefficients for the hub and mean radius would then be computed using the resulting  $\frac{1}{\sqrt{2}}$  in the polynomials, for these radii, of the form  $\frac{1}{\sqrt{2}} = \frac{1}{\sqrt{2}} \left(\frac{1}{\sqrt{2}}\right)$ . With the assumed linear variation  $\frac{1}{\sqrt{2}}$  would be

$$\mathcal{L}_{R}(2) = \mathcal{L}_{R}(1) + \left[ \frac{|r_{1}(2) - r_{1}(1)|}{r_{1}^{**} - r_{1}(1)} \right] \mathcal{L}_{R}^{**} - \mathcal{L}_{R}(1)$$
(28)

The reason for using the superscript \*\* on the mean streamline values is the same as previously mentioned in connection with Eq. 27.

For the first approximation, the Mach number  ${\rm M}_{\rm O}$  ahead of the stator is assumed, and the static properties and flowrate at station 0 are found by

$$T_{o} = \frac{T_{t_{o}}}{1 + \frac{y_{o} - 1}{2} M_{o}^{2}}$$
 (29)

$$V_o = \sqrt{rgRT_o} M_o$$
 (30)

$$\sum_{i,j} \frac{P_{i,j}}{\left(1 + \frac{K \cdot 1}{2} |M_0|^2\right)^{4/4 - 1}} \tag{31}$$

$$\beta = \frac{\beta_{-}}{\beta_{-}}$$
 (32)

$$A_{2} = \Upsilon \left( C_{1} - \Gamma_{1} \Gamma^{2} \right) \tag{33}$$

$$\dot{N} = \rho_0 A_0 V_0$$
 (34)

$$\oint A = \frac{\dot{W}}{R} \sqrt{\frac{RTr_0}{g}} = Wref$$
 (35)

The reference flowrate  $W_{\mbox{ref}}$  will be used to check continuity at stations 1 and 2.

The next step is to determine axial velocities after the stator that satisfy the equation of motion. Total enthalpy after the stator is constant by assumption, and streamline curvature effects are neglected at this stage of the analysis. With these conditions, Eq. 8 simplifies to

$$\frac{d(\ln Y_i^2)}{dX_i} = 2 TAN \, \alpha_i \frac{d\alpha_i}{dX_i} - \frac{2}{X_i} SIN^2 \, \alpha_i + \left(1 - \frac{C_i H \cos^2 \alpha_i}{Y_i^2 V_{A_i m}^2}\right) \frac{dS_i^*}{dX_i} = I$$
(36)

This equation can be integrated to give

$$/n Y_{i}^{2} = \int_{X_{0}}^{X} IdX + /n C^{2}$$

$$(37)$$

where  $X_0$  is arbitrary and  $/nC^2$  is the constant of integration. Using the boundary condition Y = 1.0 at X = 1.0,  $/nC^2$  can be found by

$$O = \int I dX + /nC^2 \quad oR \quad /nC^2 = - \int I dX$$
then
$$I \quad V^2 \quad C^2 \quad V \quad C^2 \quad V$$

 $\ln Y_{\chi}^{2} = \int I d\chi - \int I d\chi = \int I d\chi$ (38)

Equation 38 must be expressed in a form that can be utilized in the computer. Expansion by infinite series yields

$$Y = e^{\frac{1}{2} \int_{1}^{x} dX} = 1 + n + \frac{n^{2}}{2} + \frac{n^{3}}{6} + \frac{n^{4}}{24} + \frac{n^{5}}{120} + \cdots; n = \frac{1}{2} \int_{1}^{x} IdX$$
(39)

The quantities contained in I, Eq. 36, must be evaluated before proceeding with the solution. For the first approximation, assumptions are made for the values of  $Y_1$ ,  $M_1$ , and  $V_A$ . The flow angles  $\mathcal{Q}_I$  are then calculated by using Eqs. 26 and  $^m27$ . After the value of  $\mathcal{Q}_I$ , has been determined for each streamline,  $\frac{d\mathcal{Q}_I}{dX_I}$  is computed. This derivative and all others needed in the analysis are found by finite difference methods.

Enthalpy is computed by

$$H = c_{\rho} T_{t_{0}} \tag{40}$$

and the entropy term is found by the method of Vavra. 9

<sup>9</sup> Vavra, M. H., Aero-Thermodynamics and Flow in Turbomachines. New York, London: John Wiley and Sons, Inc., 1960. pp. 445-447.

$$S^* = \ln \left[ \frac{1 - \frac{Y_1^2 V_{A_1 m}^2}{C_1 H \cos^2 \alpha_1}}{1 - \frac{Y_1^2 V_{A_1 m}^2}{C_1 H \cos^2 \alpha_1 (1 - \frac{\omega_1}{2})}} \right]$$
(41)

The stator loss coefficients  $\mathcal{Y}_{S}$  are determined by equations similar to Eq. 28.

Each time a solution to Eq. 36 is found, the new values of  $Y_1$  are then used for the next iteration. After five iterations,  $\alpha_1$  and  $\frac{d\alpha_1}{d\lambda_1}$  are recalculated using the new stator exit Mach numbers. The Mach number for each streamline is found by

$$V_{A_1} = Y_1 V_{A_1 m} \tag{42}$$

$$V_{1} = \frac{V_{A_{1}}}{\cos \alpha_{1}} \tag{42a}$$

$$T_{i} = T_{t_{0}} - \frac{V_{i}^{2}}{29 \text{Jcp}}$$

$$\tag{43}$$

$$M_1 = \frac{V_1}{\sqrt{YgRT_1}}$$
 (44)

With the new values for  $\alpha_1$  and  $\frac{d\alpha_1}{dX_1}$ , five more iterations are made to determine the corrected values of  $Y_1$ .

The quantities represented by Eqs. 42-44 are recomputed after satisfying the equation of motion and additional quantities determined by

$$V_{U_1} = V_{A_1} TAN Q_1 \tag{45}$$

$$T_{1is} = T_{to} - \frac{T_{to} - T_{i}}{1 - \Psi_{s}}$$
 (46)

$$P_{i} = P_{t_{0}} \left( \frac{T_{i(s)}}{T_{t_{0}}} \right)^{\xi/\gamma - 1}$$

$$(47)$$

The flowrate through the stator is computed next and compared with that required to satisfy continuity. Reference flowrate between the hub and each streamline is found from

$$SUM = a_{n} Z r_{m} \int_{X_{H}}^{X_{a}} \xi \int dX = a_{n} Z r_{m} \int_{X_{H}}^{X_{W}} VI dX$$
 (48)

where

$$\oint = \left\{ \frac{27}{7-1} \left[ \left( \frac{P}{P_{to}} \right)^{2/7} - \left( \frac{P}{P_{to}} \right)^{5/1/7} \right] \right\}^{\frac{1}{2}}$$
(49)

The "a" and "a " in Eq. 48 are found from the polynomial which represents throat opening as a function of radius.

Before Eq. 48 is solved, the pressure ratio is compared with the critical pressure ratio. If the critical pressure ratio has been exceeded, the flow is choked at that radial location and the critical pressure ratio is used in calculating  $\phi$  and

The fraction of the total flowrate passing between the hub and each streamline is computed by

$$W_f = \frac{\sqrt{W_1 dX}}{\sqrt{W_1 dX}}$$
 (50)

Total flowrate, as found by the denominator of Eq. 50 multiplied by  $a_m \not\cong r_{mn}$ , is then compared with the reference flowrate. Overall continuity is satisfied if the difference is less than 0.0002.

If required and computed flowrates are not within this tolerance, the axial velocity for the mean streamline is adjusted by

$$V_{rm (NEW)} = V_{Alm (NEW)} + \frac{\sqrt{re_5} - W_{computed}}{0.00065}$$
 (51)

 $\mathbb{V}_{\text{req}}$  is the required reference flowrate divided by  $a_{m}\mathbb{Z}r_{m}$ , and  $\mathbb{V}_{\text{computed}}$  is the denominator of Eq. 50.

Solutions to the equation of motion and continuity are successively found until overall continuity is satisfied.

The fractions of the total flowrate determined for each streamline by Lq. 50 are then compared with the corresponding flowrate fractions ahead of the stator. Streamline continuity is satisfied if agreement is within 0.002 for each streamline. If streamline continuity has not been satisfied, the streamlines in error are adjusted by

$$X_{\text{NTN}} = X_{\text{OLD}} + \left[ W_{\text{freq}} - W_{\text{fcompreted}} \right] \frac{dX}{dW_{\text{f}}}$$
 (52)

Equation 52 applies only to streamlines 2, 3, and 4 since  $W_f = 0$  at the hub and  $W_f = 1.0$  at the tip.

If streamline positions have been adjusted, new streamline radii are found by

$$r_{NEV} = \chi_{NEW} r_{molo}$$
 (53)

and new values of X for each streamline are obtained by

$$\chi = \frac{\gamma_{\text{NEV}}}{\gamma_{\text{minimu}}} \tag{54}$$

With the new streamline radii the equation of motion and overall continuity must be satisfied again.

The following values are determined after streamline continuity is satisfied:

$$U_{i} = \frac{\text{TRFM}}{(30)(12)} r_{i}$$
 (55)

$$U_2 = \frac{r_2}{r_1^2} U_1 \tag{56}$$

$$W_{0_1} = V_{0_1} - U_1$$
 (57)

$$\beta_{1} = TAN^{-1} \left( \frac{W_{0_{1}}}{V_{A_{1}}} \right)$$
 (58)

$$W_1 = \frac{V_{A_1}}{\cos \mathcal{L}_1} \tag{59}$$

$$T_{t_{E}} = T_{i} + \frac{V_{i}^{2}}{29Jc_{P}} + \frac{U_{2}^{2} - U_{i}^{2}}{29Jc_{P}}$$
(60)

$$H_{E} = T_{t_{E}} C_{P} \tag{61}$$

$$P_{t_{E}} = P_{I} \left( \frac{T_{t_{E}}}{T_{I}} \right)^{8/\gamma - 1}$$
 (62)

Rotor blade inlet angles are determined for each streamline by means of the parabola which establishes  $\mathcal{L}_o$  as a function of  $X_1$ . Incidence is found by

$$i = \beta_1 - \beta_0 \tag{63}$$

Stalling incidences and rotor loss coefficients are then determined using the procdure described in connection with Eq. 28. If the incidence ratio  $\frac{l}{l_S}$  is less than -2.0 or greater than 1.6, the value for  $\frac{l}{l_S}$  is set equal to -2.0 or 1.6, respectively.

One of the quantities necessary for the equation of motion after the rotor is  $\frac{d \, 5 \frac{z}{z}}{d \, X_2}$ . This quantity is separated into two parts,

$$\frac{ds_2^*}{dx} = \frac{ds_{10}^*}{dx_{10}} + \frac{ds_{21}^*}{dx_{22}}$$

$$(64)$$

represents the entropy gradient due to changes of entropy across the stator and referred to station 2. The term  $\frac{c|s_{21}^*|}{c|x_2|}$  represents the entropy gradient due to the different entropy changes through the rotor. The entropy increase through the rotor is computed by using the corresponding values of the rotor in Eq. 41; for example,  $H_p$  would be used instead of H.

Neglecting streamline curvature and slope, Eq. 9 can be rewritten

$$\frac{d(\mathcal{L}_{1}Y_{1}^{2})}{dX_{2}} = -i TAN \beta_{z} \frac{d\mathcal{B}_{z}}{dX_{2}} SIN^{2} \beta_{z} - \frac{4 Um \cos \beta_{z} SIN \beta_{z}}{Y_{2} WA_{2m}}$$

$$-\frac{2U_mU_2\cos^2\beta_2}{Y_2^2W_{A_{2m}}^2} + \frac{C_1\cos^2\beta_2}{Y_2^2W_{A_{2m}}^2} \frac{dH_E}{dX_2} + \left(1 - \frac{C_1H_E\cos^2\beta_2}{Y_2^2W_{A_{2m}}^2}\right) \frac{dS_2^*}{dX_2}. \tag{65}$$
The discharge angle  $\beta_2$  and its derivative  $\frac{d\beta_2}{dX_2}$  are found in the same manner as previously described for the stator discharge angles. The method for solving Eq. 65 is similar to that used for Eq. 36 with the exception that iterations are carried out until correspond-

with the exception that iterations are carried out until corresponding values of Y change by less than 0.005, or until 13 iterations have been completed. The extra iterations are necessary because there may be a slower convergence of the values of Y at station 2. After satisfying the equation of motion, values at station 2 corresponding to those represented by Eqs. 42-47, are computed. It should be noted that where absolute velocity terms are used in Eqs. 42-47, the corresponding equations for the rotor will employ relative velocities. Therefore M<sub>2</sub> is the Mach number of the flow relative to the rotating rotor blade.

Overall continuity is checked at the rotor discharge using the same procedure utilized for the stator; however, the reference flowrate is increased according to Eq. 25 to account for the flowrate between the blade tips and the surrounding shroud.

After overall continuity has been satisfied at station 2, streamline continuity is checked. In addition to the functions performed for the check after the stator, there are certain quantities that must be adjusted for streamline relocation. These are

$$\left(\frac{dH_{E}}{dX_{2}}\right)_{NEW} = \left(\frac{dH_{E}}{dX_{2}}\right)_{OLD} + \frac{dX_{2}}{dX_{2}}\left(X_{NEW} - X_{OLD}\right)$$
(66)

$$H_{L_{NEV}} = H_{\Gamma_{OLL}} + \frac{JH_{E}}{JX_{2}} (X_{NEW} - X_{OLD})$$
 (67)

$$\left(\frac{J s_{10}^{*}}{d N_{2}}\right)_{NFW} = \left(\frac{d s_{10}^{*}}{d N_{2}}\right)_{OLD} + \frac{d^{2} s_{10}^{*}}{d N_{2}^{2}} \left(\chi_{NFW} - \chi_{OLD}\right)$$
(68)

$$U_{2,NEN} = \left(U_{2,OLD}\right) \frac{\chi_{NEW}}{\chi_{2,DD}} \tag{69}$$

A solution exists when the equations of motion, and overall and streamline continuity have been satisfied; however, this solution has neglected streamline curvature and streamline slope. To obtain a solution which accounts for streamline curvature effects, the terms in Eqs. 8 and 9 that contain Sr or  $\Delta R$  must be included when solving the equations of motion.

The streamline displacements  $\leq r$  and  $\Delta R$ , shown in Fig. 2. are computed for each streamline by

$$Sr = r_1 - \frac{r_0 + r_2}{2} \tag{70}$$

and

$$\Delta E = V_0 - V_2 \tag{71}$$

Equation 70 is based on the assumption that the radii of any streamline at stations 0 and 2 are not greatly different. Therefore the cosine of the angle between the line connecting these points and the axis of the machine is approximately equal to unity. The length L, also shown in Fig. 2, is taken to be half the distance from 0.1 inch ahead of the stator to 0.1 inch after the rotor. After a solution is obtained which accounts for streamline curvature effects, the resultant pressure ratio  $\frac{\rho_{to}}{\rho_2}$  is compared with the  $\frac{\rho_{to}}{\rho_2}$  specified initially. Ideally  $\frac{\rho_{to}}{\rho_2}$  for each streamline will be the same. However, computer solutions for this quantity may vary slightly from streamline to streamline. Therefore a mass-flow-weighted value of  $\frac{\rho_{to}}{\rho_2}$ , called  $\left(\frac{\rho_{to}}{\rho_2}\right)_{\dot{W}}$ , is found by

$$\left(\frac{P_{t_o}}{P_2}\right)_{\dot{W}} = \sum_{i=1}^{4} \left(\frac{P_{t_o}}{P_2}\right)_{(i+1)} + \frac{P_{t_o}}{P_2}\left(\frac{W_{f(i+1)} - W_{fi}}{2}\right)$$
(72)

If the specified  $\frac{P_{to}}{P_2}$  and  $\left(\frac{P_{to}}{P_2}\right)$ , differ by more than 0.0003, the Mach number of the flow ahead of the stator is properly adjusted and another solution is found.

After the iterations for  $\frac{P_{to}}{P_{z}}$  have been completed, additional quantities are determined for each streamline, using

$$\Delta H = H_1 - H_2 = (U_1 V_{U_1} - U_2 V_{U_2}) \frac{1}{9J}$$
 (73)

$$T_{tz} = T_{to} - \frac{\Delta H}{C_P} \tag{74}$$

$$T_{2is} = T_{t_o} \left( \frac{P_z}{P_{t_o}} \right)^{\frac{N-1}{N}}$$
 (75)

Overall efficiency is then computed by

$$\gamma = \frac{\Delta H}{\Delta h_{is}} = \frac{T_{to} - T_{tz}}{T_{to} - T_{zis}}$$
The ideal change in enthalpy  $\Delta h_{is}$  is the isentropic enthalpy (76)

The ideal change in enthalpy  $\Delta h_{is}$  is the isentropic enthalpy drop from the total inlet pressure  $P_{to}$  to the static discharge pressure  $P_{2}$ . Equation 76 is used to compute the efficiency of a

single stage turbine because the kinetic energy leaving the rotor cannot be utilized. The efficiency defined by En. 76 will be referred to as total-static efficiency.

Theoretical degree of reaction  $r^*$  and head coefficient  $k_{\mbox{is}}$  are given by

$$r^* = \frac{T_{i\dot{c}s} - T_{z\dot{c}s}}{T_{\bar{c}o} - T_{z\dot{c}s}} \tag{77}$$

and

$$k_{is} = \frac{\Delta h_{is}}{U_i^2/2gJ} = \frac{2C_P(T_{to} - T_{zis})gJ}{U_i^2}$$
(78)

For turbine performance curves it is desirable to obtain the mass-flow-weighted values of efficiency ( ) $_{\dot{w}}$ , head coefficient (k $_{\dot{i}\dot{s}}$ ) $_{\dot{w}}$ , theoretical degree of reaction (r\*) $_{\dot{w}}$ , horsepower (HP) $_{\dot{w}}$ , and moment (M $_{\rm p}$ ) $_{\dot{w}}$ . The last two quantities are found from

$$(HP)_{\dot{w}} = \frac{(\Delta H)_{\dot{w}} J_{\dot{w}}}{550} \tag{79}$$

and

$$(M_R)_{\dot{W}} = \frac{(HP)_{\dot{W}}(550)}{\omega}$$

The mass-flow-weighted  $\Delta H$ , called  $(\Delta H)_{\dot{w}}$ , as well as  $(^{\dot{\gamma}})_{\dot{w}}$ ,  $(^{\dot{k}}_{\dot{i}\dot{s}})_{\dot{w}}$ , and  $(^{r*})_{\dot{w}}$  are computed using equations similar to Eq. 72.

Referred values are obtained, following NASA practice. For  $\gamma$  = 1.4:

$$HP_{ref} = \frac{(HP)_{\dot{W}}}{\sqrt{\Theta} \delta}$$
 (81)

$$M_{R \text{ ref}} = \frac{(M_R)\dot{w}}{6}$$
 (82)

$$RPM_{fef} = \frac{RPM}{V\Theta}$$
 (83)

$$\dot{W}_{ref} = \frac{\dot{W} V \Theta}{\delta}$$
 (84)

$$V_{YEF} = \frac{V}{\sqrt{Q}}$$
 (84a)

where:

$$\theta = \frac{T_{to}}{T_{57D}} = \frac{T_{to}}{5/8.4} \tag{85}$$

$$S = \frac{P_{to}}{P_{STO.}} = \frac{P_{to}}{14.7} \tag{86}$$

## 4. MOD I and MOD II Turbines

The method of analysis as presented was used to determine performance curves for the so-called MOD I and MOD II turbines. Both turbines are single stage axial-flow machines. Experimental tests were conducted on the MOD II turbine by Commons and Messegee and are described in Refs. 4 and 6. The test results are plotted with appropriate predicted performance curves.

The so-called MOD I turbine was designed for free-vortex flow and has highly twisted blades. Outer diameter of this turbine is 9.898 inches. The hub diameters of the stator inlet and rotor discharge are 6.930 and 5.970 inches, respectively. The stator contains 13 blades and the rotor 22 blades. Blades of the MOD I turbine are generally thin. The stator and rotor profiles used to predict outlet angles and loss coefficients are shown in Fig. 4.

The MOD II turbine is approximately the same size as the MOD I, but its blading is considerably different. The blades of the MOD II turbine are thick with blunt leading edges and constant profiles over the blade height. Outer diameters of the MOD II turbine stator and rotor are 9.701 and 9.836 inches, respectively. The stator has a hub diameter of 6.796 inches, and the hub diameter of the rotor is 6.598 inches. There are 19 stator blades and 18 rotor blades. Stator and rotor blade profiles for the MOD II turbine are shown in Fig. 5.

Throughout the remainder of this thesis the MOD I and MOD II turbines will be referred to simply as MOD I and MOD II.

The minimum throat opening "a" of the blade channels is a very critical dimension. Slight variations of this quantity have a considerable effect on turbine performance. Since this quantity is so sensitive, values of "a" measured from the actual hardware were used for the analysis rather than those obtained from the manufacturing

drawings. Then errors due to manufacturing will not be a factor in comparison of predicted and experimental results. Figure 6 shows the throat openings "a" as a function of radius for the stator and rotor blading of both turbines.

Predicted stator outlet angle  $\alpha_1$  as a function of radius  $r_1$  for the MOD I is plotted in Fig. 7. The assumed linear variation of  $\alpha_1$  between the hub and the mean radius and between the mean and the tip is readily apparent in this figure. Figure 8 shows the predicted variations of  $\alpha_1$  with Mach number  $\alpha_1$  for the MOD I. Figures 9 and 10 are the corresponding plots for the MOD II.

Relative discharge angles  $\mathcal{B}_2$  were computed for two radial tip clearances k for each turbine; namely, 0.020 and 0.033 inches for the MOD I, and 0.015 and 0.033 inches for the MOD II. The predicted flow angles  $\mathcal{E}_1$  are plotted in the same manner as previously described for the flow angles  $\alpha_1$ . Figures 11 and 12 show  $\mathcal{B}_2$  as a function of radius  $\alpha_2$  and as a function of Mach number  $\alpha_2$ , respectively, for the MOD I, where  $\alpha_2$  refers to the Mach number of the flow relative to the rotor. Figures 13 and 14 show the corresponding plots for the MOD II. The predicted effect of radial tip clearance on the discharge angles  $\alpha_2$  can be seen in Figs. 11-14.

Stator loss coefficients  $\mathcal{C}_{S}$  were computed for the radial locations corresponding to the hub, mean radius, and tip. Figure 15 shows  $\mathcal{C}_{S}$  as a function of radius  $r_{1}$  for both the MOD I and MOD II. The straight lines in this figure between the values of at the hub and mean radius, and between the mean and the tip, reflect the assumed linear variation of  $\mathcal{C}_{S}$  in these regions.

Variation of rotor blade inlet angle  $\mathcal{L}_c$  with radius  $r_1$  is plotted for both turbines in Fig. 16. The difference between the untwisted and the free-vortex blades is easily noted in this plot. Also shown in this figure are curves representing the variation of stalling incidence  $i_s$  with  $r_1$ . The change of the MOD I rotor blade profiles with radius is reflected by a considerable variation of  $i_s$  whereas just the opposite is true for the MOD II.

Curves for the MOD I showing predicted rotor loss coefficients as a function of incidence ratio  $\frac{i}{i_s}$  for the hub, mean radius, and tip are shown in Fig. 17. Since the loss coefficients between the mean radius and the tip are dependent on tip clearance, there are two curves for the tip. One curve holds for the tip clearance of 0.020 inches; the other is for the larger tip clearance of 0.033 inches. For negative incidence ratios between -0.5 and -2.0 the curves for the tip are estimations. This was necessary because the computations by the method shown in Appendix B, Section 2, gave unrealistically low values of  $\mathcal{J}_R$  as the flow inlet angle  $\mathcal{L}_i$  approached -90°. The situation is more easily understood when it is noted that the blade inlet angle is -33.7° and the stalling incidence is 37° at the rotor blade tip. Figure 17 shows that the loss coefficients for the hub  $(r_1=3.300 in.)$  are relatively large. The loss coefficients at the hub are larger for most incidence ratios than those at the tip for a tip clearance of 0.020 inches. The larger value of at the hub reflects the higher losses that are associated with an impulse type blade. It can be seen in Fig. 4 that the blade shape varies from a reaction type profile at the tip to an impulse type profile at the hub.

Loss coefficients for the MOD II rotor are plotted in Fig. 18 for tip clearances of 0.015 and 0.033 inches. The predicted similarity of the curves for the hub and mean radius is to be expected since the blading differs only in solidity. Although the blade profile is the same at all radii, the losses due to tip clearance result in larger loss coefficients for the tip profile.

For convenience, the blade properties used for calculating the MOD II rotor loss coefficients were those at the hub, mean and tip radii of the rotor discharge. Since the annulus area at the rotor discharge is larger than the annulus area at the stator discharge and since the flow incidence is a significant parameter for the rotor loss coefficients, it would have been more appropriate to use the blade characteristics at the hub, mean, and tip radii of the rotor inlet. However, the error is insignificant because the blade profile does not change along the radius.

It may be noted from Fig. 4 that the minimum radius, for which a blade profile is given, is 3.597 inches whereas the radius at the hub of the MOD I stator exit is 3.300 inches. The outlet angle and loss coefficient for the hub were found by extrapolation, using the values computed for the radii of 4.125 and 3.597 inches and assuming a linear variation of these quantities with radius. The relative outlet angle for the hub of the MOD I rotor was found in a similar manner.

Performance curves for the two turbines were determined for the rotor tip clearances mentioned earlier. The total inlet to static discharge pressure ratios investigated were 1.30, 1.40, 1.50, and 1.60, with the exception that pressure ratios of 1.31 and 1.51 were used for the MOD II with 0.033 inch rotor tip clearance. These pressure ratios agree more closely with those experimentally investigated by Commons and Messegee.

The  $\epsilon$  xial distances L used for the determination of the curvatures depend on the axial clearance between the stator and rotor as well as on the blade geometries. Axial clearances of 0.4 and 1.0 inches were used for the analyses of the MOD I and MOD II, respectively.

Curves representing the performance of the MOD I are plotted in Figs. 19 through 26. Performance values plotted are mass-flow-weighted values unless stated otherwise. Figure 19 shows referred flowrate as a function of referred RPM. The increase in flowrate due to an increase in rotor tip clearance can be seen in this figure. Although the flowrate is greater for the larger tip clearance, the torque developed is greater at the smaller tip clearance. The decrease in torque for the larger tip clearance results from the increased losses and decrease in turning angle of the flow through the rotor near the tip. The predicted effect of tip clearance on torque is shown in Fig. 20 where referred moment is plotted versus referred RPM.

The variation of total-static efficiency with referred RPM for the two tip clearances can be seen in Figs. 21 and 22. The referred RPM at which maximum efficiency occurs increases when the total inlet to static discharge pressure ratio is increased. Blade losses and the kinetic energy of the flow leaving the rotor affect the total-static efficiency. As pressure ratio is increased the absolute velocity leaving the stator and the relative velocity leaving the rotor increase. Therefore the peripheral speed of the rotor must increase to obtain conditions where the absolute velocity leaving the rotor is in an axial direction and where the relative flow ahead of the rotor has zero incidence. The RPM where the flow has zero incidence on the rotor will not necessarily be that at which the absolute velocity leaving the rotor is axial. At any RPM, flow incidence and absolute discharge angles vary from streamline to streamline, and the above statements refer to mass-flow-weighted values.

It may be noted in Figs. 21 and 22 that the peak total-static efficiency decreases somewhat as the pressure ratio increases. At higher pressure ratios the ratio of kinetic energy leaving the rotor to the work done on the rotor increases. Since the kinetic energy leaving the rotor is lost energy for a single stage turbine, the total-static efficiency declines. The effects of pressure ratio and tip clearance on efficiency can be seen in Fig. 23 where total-static efficiency is plotted as a function of the isentropic head coefficient.

The variation of referred power with referred RPM can be seen in Fig. 24. Peak power does not occur at the same referred RPM at which peak efficiency occurs. The peak referred power occurs at the referred RPM where the product of total-static efficiency and referred flowrate is greatest.

Theoretical degree of reaction is plotted as a function of isentropic head coefficient in Fig. 25. It may be noted that the theoretical degree of reaction increases with increasing pressure ratio and decreases with increasing tip clearance for any given isentropic head coefficient. The predicted effects are considerably different from the results of the radial turbine tests conducted by Riley. 10 Riley found that theoretical degree of reaction was independent of pressure ratio and axial clearance for radial turbines.

<sup>10</sup> Riley, M. W., The Effect of Axial Clearance on the Performance of a Dual Discharge Radial Turbine (USNPG Thesis, December 1966), p. 70.

Performance values for each streamline are obtained from the computer solution. However, plots utilizing values for each streamline would be difficult to analyze. The deviation of the hub and tip values from that of the mass-flow-weighted average may be seen in Fig. 26. In this figure, hub, tip, and mass-flow-weighted values of theoretical degree of reaction are plotted as functions of referred RPM for a total inlet to static discharge pressure ratio of 1.40.

Performance curves for the MOD II corresponding to those presented for the MOD I are plotted in Figs. 27 through 39. Additional plots have been used for the MOD II because of the inclusion of experimental results. An axial clearance of 1.0 inches was used for the theoretical prediction. Therefore, only experimental data for that axial clearance are shown.

Comments made concerning the performance curves of the MOD I are applicable to the MOD II performance curves also. Differences in the performance of the two turbines will be discussed later.

Plots of the variation of referred flowrate with referred RPM are shown in Figs. 27 and 28 for two rotor tip clearances. The quantitative values as well as the curve shapes agree well with the experimental data. The maximum difference between predicted and experimental referred flowrates occurs at a pressure ratio of 1.51 for a tip clearance of 0.033 inches. There are two experimental points for this pressure ratio and tip clearance that differ from the predicted curve by about 2 per cent. The predicted and experimental values for all pressure ratios and tip clearances have an average difference of less than 1 per cent.

Figures 29 and 30 show curves of referred moment versus referred RPM. The trends expressed by the predicted curves are in excellent agreement with the experimental data. Although the quantitative agreement between theoretical and experimental values is very good for three of the curves, the experimental torque is generally lower than the predicted torque. The average difference between predicted and experimental values is about 3 per cent.

Figures 31 through 34 show total-static efficiency as a function of referred RPM. The shapes of the predicted curves generally agree well with the experimental results. However, there is an indication that experimental efficiencies decrease more rapidly at high RPM than is predicted by the theoretical curves. In the high RPM region the upper part of the rotor blade has a large negative flow incidence. In the prediction analysis when the incidence ratio had a value less than -2.0, the value of -2.0 was used for computing loss coefficients. This limitation may be the reason that the predicted efficiencies in the high RPM range do not decrease as rapidly as the test data indicate.

The quantitative agreement between predicted and experimental efficiencies varies considerably between different pressure ratios and tip clearances. There are two data points at a pressure ratio of 1.31 and tip clearance of 0.033 inches where the experimental efficiencies are over five points below the predicted values. At a pressure ratio of 1.50 and a tip clearance of 0.015 inches the average difference between predicted and experimental efficiencies is 1.5 points. Giving equal weight to all experimental values the average difference between experimental and predicted values is 2.6 points.

The calculations gave a decrease in efficiency by about two points as the tip clearance was increased from 0.015 to 0.033 inches. The decrease in experimental efficiencies for the increased tip clearance varied with the different pressure ratios. However, the average decrease in efficiency is close to the predicted decrease.

Total-static efficiency as a function of isentropic head coefficient  $k_{is}$  is shown in Fig. 35 for pressure ratios of 1.40 and 1.60. Figure 38 shows degree of reaction r\* versus  $k_{is}$ . No experimental data are plotted in these figures because experimental values of mass-flow-weighted r\* and  $k_{is}$  were not available.

Plots of referred power as a function of referred RPM are shown in Figs. 36 and 37. The comments made earlier concerning the referred moment plots apply to these curves also, since the turbine power is proportional to the product of torque and RPM.

Hub, tip, and mass-flow-weighted values of theoretical degree of reaction are plotted in Fig. 39 for a pressure ratio of 1.40. The experimental values for the hub and tip ere also plotted in this figure. The trend of the experimental data is in close agreement with the predicted trend; however, the quantitative agreement is poor, especially at the hub. The decrease in degree of reaction, as the tip clearance is increased, is considerably greater than predicted. Both predicted and experimental results indicate that an axial flow turbine differs from a radial turbine inasmuch as the theoretical degree of reaction changes if rotor tip clearance is changed.

Figure 40 shows the absolute flow angles  $\mathcal{N}_1$  and  $\mathcal{N}_2$  as a function of radius for the stator and rotor outlets, respectively, of the MOD II. Experimental data were obtained for the same operating conditions; namely,  $P_{t_0}/P_2=1.40,k=0.015$  inches, RPM/ $\sqrt{2}=13,394$ . The experimental and predicted stator outlet angles have an average difference of less than half a degree over the blade height. Predicted and experimental agreement for the absolute flow angles at the rotor outlet is not as close. If the experimental point near the hub is neglected, the average difference between predicted and experimental values of  $\mathcal{N}_2$  is about 3.5 degrees. The experimental flow angles  $\mathcal{N}_2$  are larger than the predicted angles over most of the blade height, which explains why the predicted torque was generally larger than the measured torque.

Velocity distributions at the stator and rotor exits of the MOD II are shown in Fig. 41 for the operating conditions listed in the preceding paragraph. The experimentally found stator discharge velocities varied only slightly more from hub to tip than predicted. The average difference between predicted and experimental stator exit velocities was about 1 1/2 per cent.

The experimentally determined absolute rotor exit velocities, which are shown in Fig. 41, are considerably larger than the predicted velocities except at the hub and tip. Although experimental points could be taken only at distances greater than 0.16 inches from the inner diameter of the shroud, the trend of the data points indicates that the experimental velocities are less than predicted at the tip.

Near the hub the experimental velocities were less than predicted. The decrease in experimental velocities near the hub and tip is the result of separated flow in these regions. The reader is referred to Ref. 6 for additional information concerning the experimental data and for photographs of the rotor showing indications of separation.

Experimental relative rotor discharge angles  $\mathcal{L}_2$  were determined for the known rotor speed from the measured values of absolute rotor outlet angles and velocities, which are plotted in Figs. 40 and 41, respectively. The experimental and predicted values of  $\mathcal{L}_1$  are shown in Fig. 42. The magnitudes of the experimental values of  $\mathcal{L}_2$  near the hub are larger than predicted. Over approximately the outer three fourths of the rotor blade height the magnitudes of the experimental values of  $\mathcal{L}_2$  are less than predicted. At a radius of 4.4 inches, for instance, the experimental value of  $\mathcal{L}_2$  is about -57 degrees whereas the predicted value is about -67 degrees. The average difference between experimental and predicted values of  $\mathcal{L}_2$  is about 6.5 degrees. The lower values of  $\mathcal{L}_2$  predicted near the hub, provide some compensation in the overall performance for the high values of  $\mathcal{L}_2$  predicted for the outer part of the blade.

The MOD II was not designed to achieve the highest possible efficiency. The objectives of the design were to investigate the effects of b unt untwisted bladings of constant profile. This type of blading has the following advantages:

- 1. Blade cooling passages are easily accommodated.
- A wider operating RPM range is possible for a specified turbine efficiency variation.
- 3. Constant profile blades can be manufactured more economically than twisted blades, especially if exotic high strength materials are needed for elevated temperatures.

The following discussion concerning comparison of the MOD I and MOD II is based entirely on the predicted performance of these turbines using the methods described in this thesis. The reader should be aware that Ainley's methods for predicting loss coefficients were not developed for use with rotor blades having blunt leading edges.

Therefore the accuracy of the rotor loss coefficients predicted for the MOD II is probably not as good as for the MOD I.

Figure 43 shows the predicted axial velocity ratios as a function of radius ratio at the stator exit and at the rotor exit for the MOD I and the MOD II. The curves are for the conditions that occur at the maximum efficiency of each turbine. The axial velocities of the MOD I are nearly constant over the blade height at the stator exit and at the rotor exit. The blading of this turbine was originally designed for free-vortex flow by assuming uniform axial velocity components from hub to tip. From Fig. 43 it can be seen that this condition is only approximately satisfied since loss variations and curvature influences produce slight variations from the assumed distribution. Axial velocities for the MOD II decrease from the hub to the tip at the stator exit, and increase from the hub to the tip at the rotor exit, since the bladings of this turbine have constant profiles. The variation of axial velocities for the MOD II is more pronounced at the rotor exit than at the stator exit. At  $\frac{\text{VA}_2}{\text{VA}_2}$  varies from 0.735 at the hub to 1.195 at the the rotor exit tip.

The maximum predicted total-static efficiency of the MOD I at a tip clearance of 0.033 inches is about 84 per cent. The MOD II has a maximum predicted efficiency at that tip clearance of about 80 per cent. The difference in predicted peak efficiencies is due to different factors. The stator loss coefficients are higher for the MOD I than for the MOD II, whereas the loss coefficients for the MOD II rotor at zero incidence are comewhat larger than those for the MOD I at zero incidence. An accurate comparison of loss coefficients is more involved for the rotor than for the stator. Because of its twisted blades, the MOD I rotor has nearly zero incidence at all points along the blade height in the vicinity of the optimum efficiency. Only at one radius of the MOD II rotor blade, however, is the incidence zero. The incidence angles at larger radii are negative and at smaller radii they are positive. Therefore part of the blade is always operating at a loss coefficient larger than the minimum. Of probably even greater significance

for the efficiency decrease is the kinetic energy of the gas leaving the turbine. Minimum kinetic energy is lost when the flow is axial in direction. The design of the MOD I is such that at all points along the blade height the absolute velocity at the rotor exit is nearly axial in the vicinity of the optimum efficiency. The MOD II has only a small radial portion of the rotor blade where the discharge angle  $\alpha_2$  is zero. At radii greater than the radius where  $\alpha_2$  is zero, the absolute flow discharge angle is positive and at smaller radii it is negative. Therefore the kinetic energy leaving the rotor of the MOD II is greater than that of the MOD I when both turbines are operating at peak efficiency.

A comparison of the total-static efficiency versus isentropic head coefficient curves for the MOD I (Fig. 23) and MOD II (Fig. 35), shows a larger variation of efficiency with head coefficient for the MOD I. For example, at k=0.033 inches and  $\frac{P_{C_0}}{P_2}$ =1.60, the MOD I efficiency decreases about 11.5 points as the head coefficient increases from 3 to 7 whereas the MOD II efficiency decreases only about 7 points for the same change in head coefficient. This difference in efficiency variation results from the effects of blade twist mentioned earlier.

To investigate the influence of the streamline curvatures on the predicted performance of the MOD I and MOD II the presented flow equations were solved for an axial length of the bladings L of 9 x 10<sup>5</sup> inches. Increasing L to this value has the same effect as neglecting streamline curvature effects as can be seen from Eqs. 124 and 125. Differences between performance values neglecting streamline curvature, and those where curvature effects were accounted for were less than 0.2 per cent for both turbines.

## 5. Conclusions and Recommendations

Effects of streamline curvature were found to be insignificant for the MOD I and MOD II. The small effect on predicted turbine performance due to streamline curvature indicates that the exact value assumed for the curvature factor K will not greatly affect predicted performance. The results of this analysis also indicate that Vavra's method of approximating streamline curvature is of sufficient accuracy for methods of analysis where the flow equations are satisfied at stations between blade rows.

It is recommended that the MOD I be tested. Results of that investigation would provide additional information concerning the general applicability and accuracy of the method of analysis proposed in this thesis.

Predicted and experimental flowrates for the MOD II were in close agreement. Since the restriction factor  $\xi$  is a significant factor in the predicted flowrates, the experimental results verify the validity of the theory used in the development of  $\xi$ .

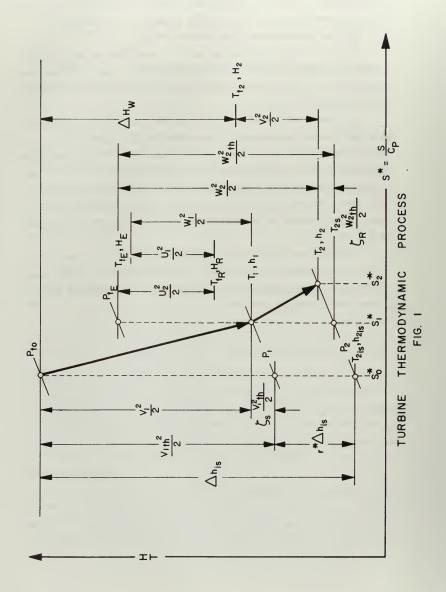
Experimental results for the MOD II showed that the predicted rotor torque and turbine efficiency were generally 2 to 3 per cent too high. The angles and velocities predicted for the flow at the stator outlet were in excellent agreement with the experimental results. Therefore the high values predicted for rotor torque and turbine efficiency must be tied to the rotor solution. The measured values of outlet angles and velocities for the rotor discharge also indicate that the predicted rotor solution is not completely correct.

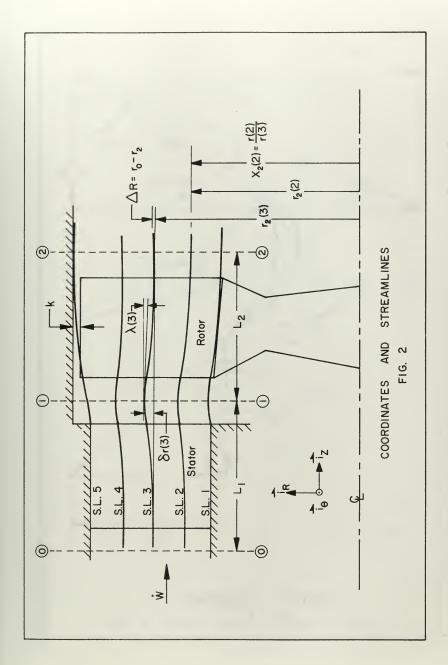
Exactly what parameters are in error in the predicted rotor solution, and to what extent, cannot be stated without more experimental data. More traverses should be taken at the stator and rotor outlets using recently calibrated probes. With this information,loss coefficients and relative discharge angles for the rotor could be determined on a streamline basis. Based on the experimental results included in this thesis,it is suspected that the magnitudes of the predicted discharge angles  $\beta_2$  are too large. If the predicted turning angles of the flow through the rotor were less, the predicted torque and efficiency would decrease. Also, the separated flow at the hub and tip at the discharge of the MOD II rotor indicates that predicted loss coefficients should be higher for the streamlines at these locations.

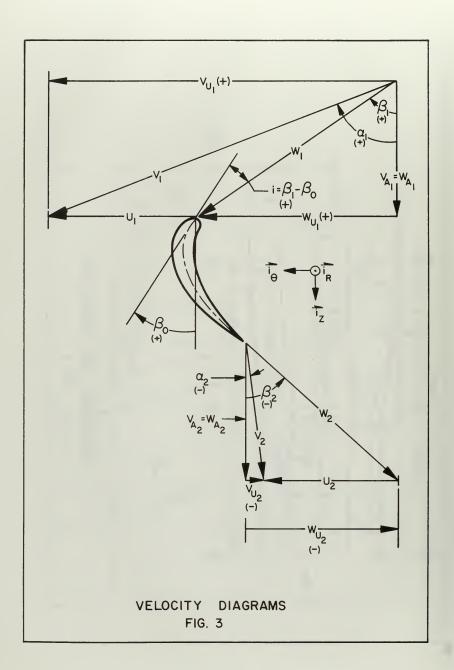
To accurately predict turbine performance at off-design as well as design conditions, a streamline analysis is necessary. The experimental verification obtained for the proposed method indicates that the method possesses much potential and that additional development

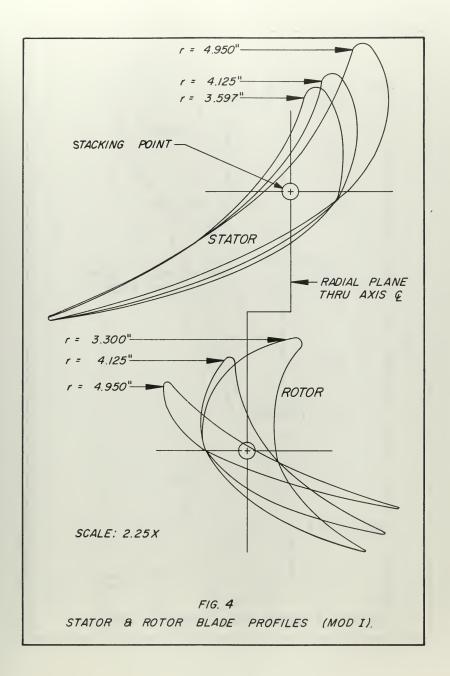
is warranted. For future development work on this type of analysis it is recommended that the following changes be considered:

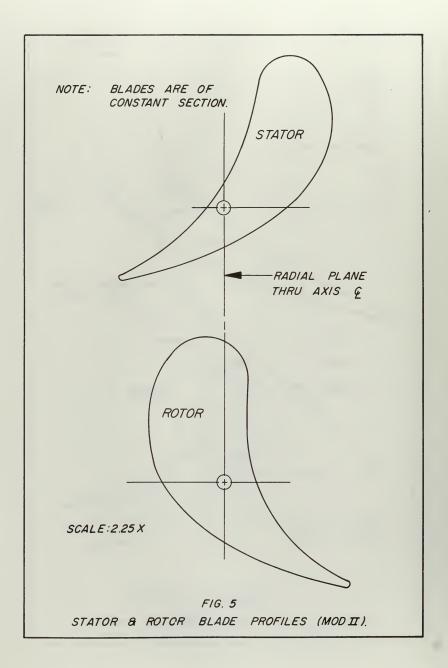
- 1. Use seven streamlines instead of five.
- 2. Calculate discharge angles using the methods of this thesis but neglecting blade curvature effects. That would decrease the angles by 2 to 3 degrees and be in closer agreement with experimental results.
- 3. Introduce a multiplying factor for the rotor loss coefficients which would more accurately account for the separated flow at the hub and tip but would not greatly change the average rotor loss coefficient. For example, rotor loss coefficients would first be calculated using the methods of this thesis, and then the loss coefficients for streamlines 1, 2, 6, and 7 would be multiplied by, say, 1.4 and those for streamlines 3, 4, and 5 would be multiplied by 0.6 or 0.7.
- 4. Calculate rotor loss coefficients for an range of -3.0 to 2.0 instead of -2.0 to 1.6. The increased range would improve performance prediction for turbines with untwisted rotor: blades at off-design conditions.

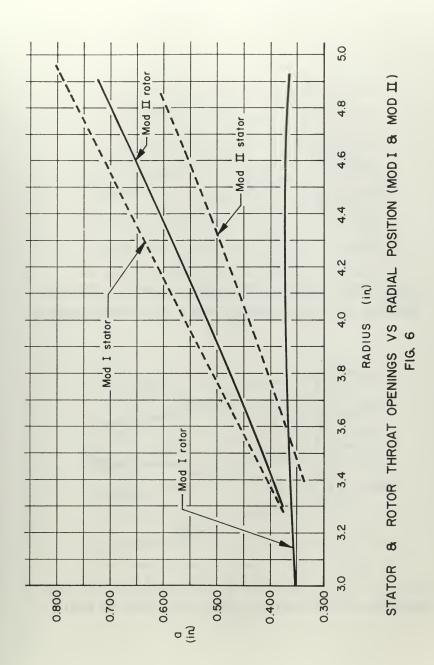


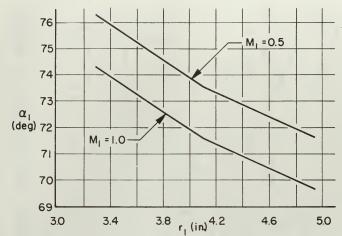




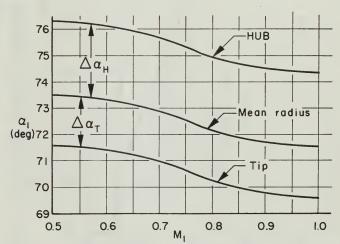




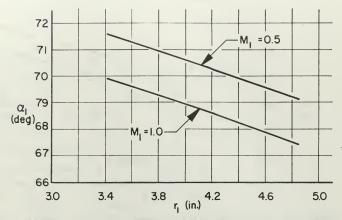




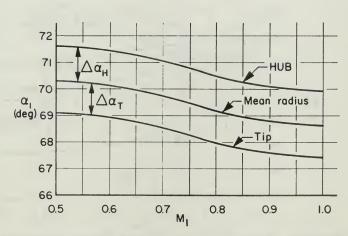
VARIATION OF STATOR OUTLET ANGLE WITH RADIUS (MOD I) FIG. 7



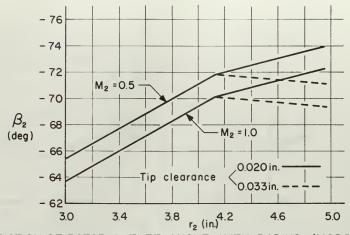
VARIATION OF STATOR OUTLET ANGLE WITH MACH NO. (MODI) FIG. 8



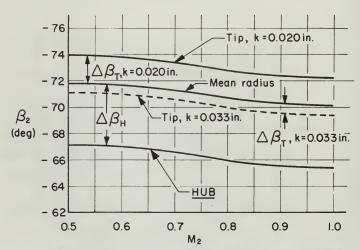
VARIATION OF STATOR OUTLET ANGLE WITH RADIUS (MOD II)
FIG. 9



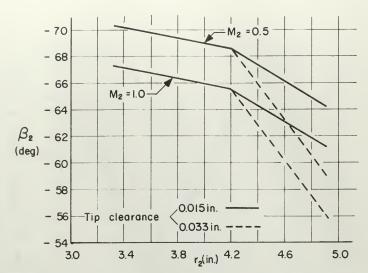
VARIATION OF STATOR OUTLET ANGLE WITH MACH NO. (MODII) FIG. 10



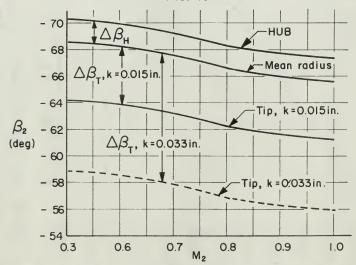
VARIATION OF ROTOR OUTLET ANGLE WITH RADIUS (MODI) FIG. 11



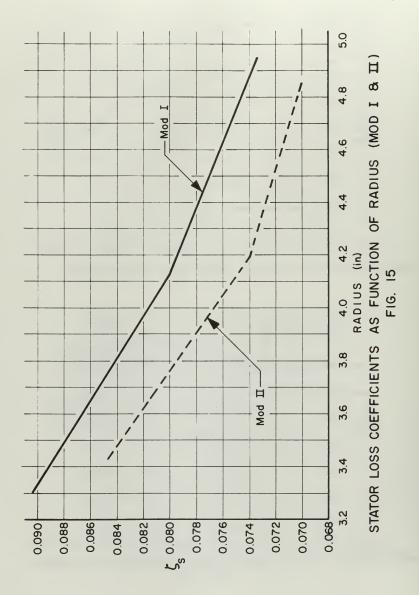
VARIATION OF ROTOR OUTLET ANGLE WITH MACH NO. (MODI) FIG. 12

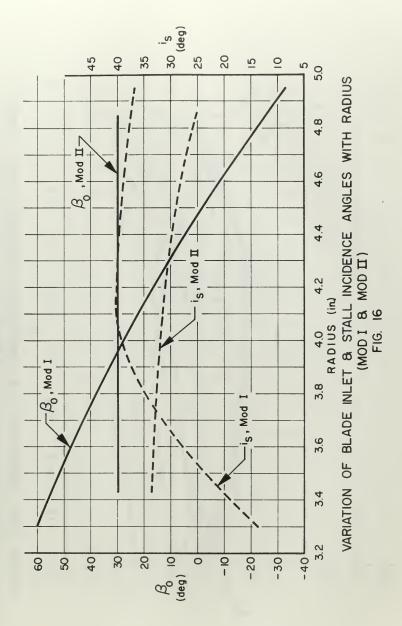


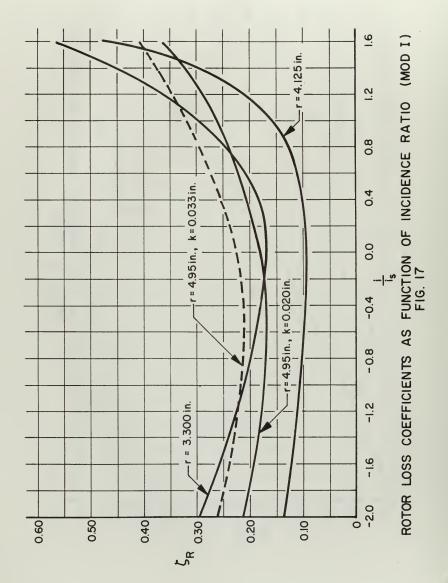
VARIATION OF ROTOR OUTLET ANGLE WITH RADIUS (MOD  $\Pi$ ) FIG. 13

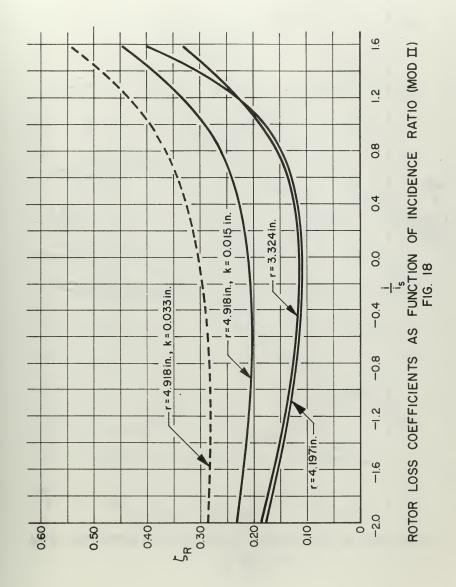


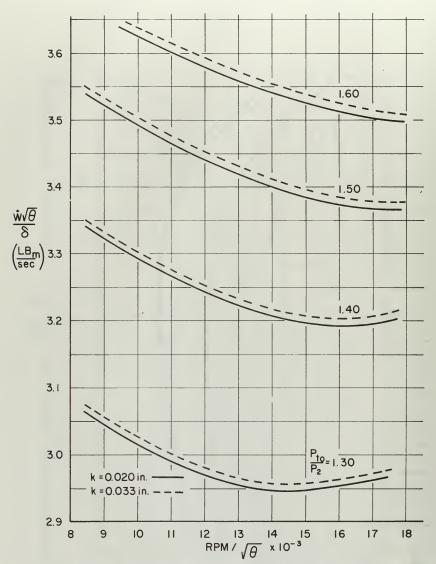
VARIATION OF ROTOR OUTLET ANGLE WITH MACH NO. (MODIL) FIG. 14



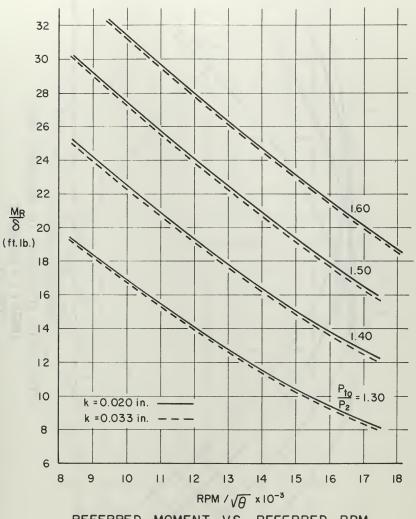




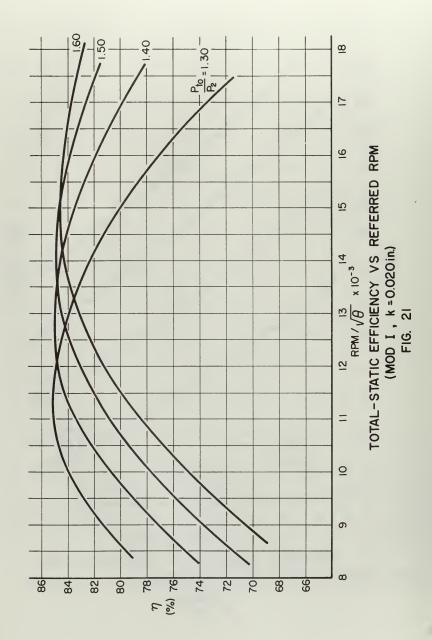


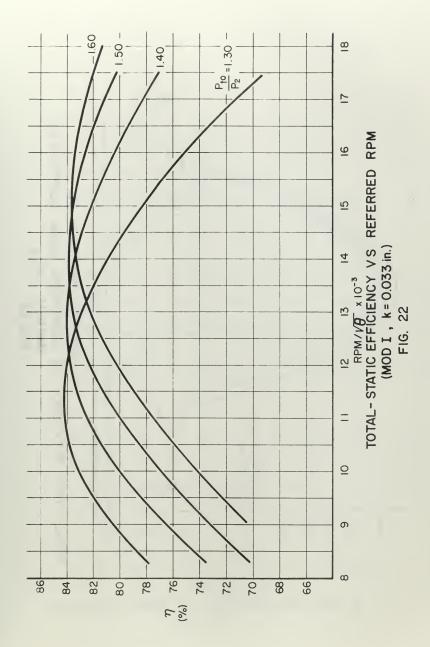


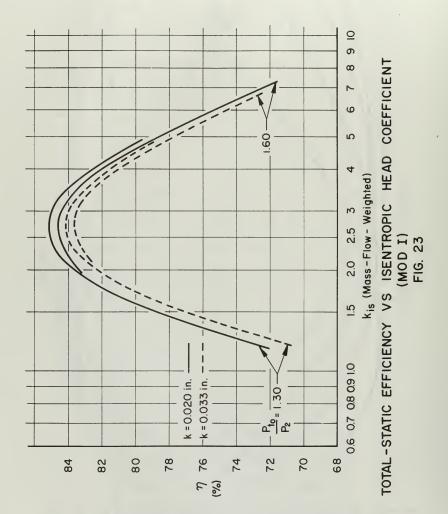
VARIATION OF REFERRED FLOWRATE WITH REFERRED RPM (MOD I)
FIG. 19



REFERRED MOMENT VS REFERRED RPM (MOD I)
FIG. 20







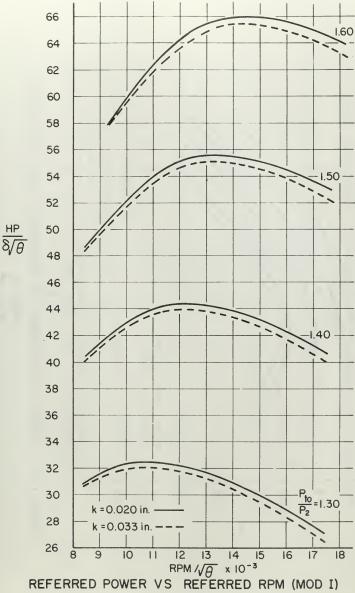
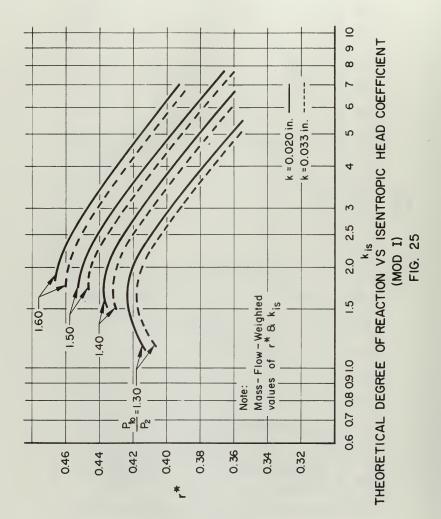
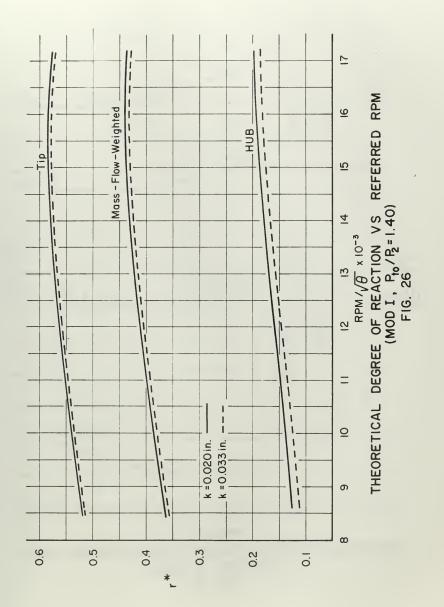
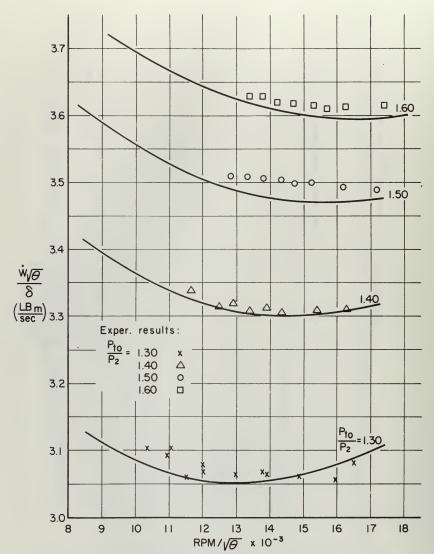


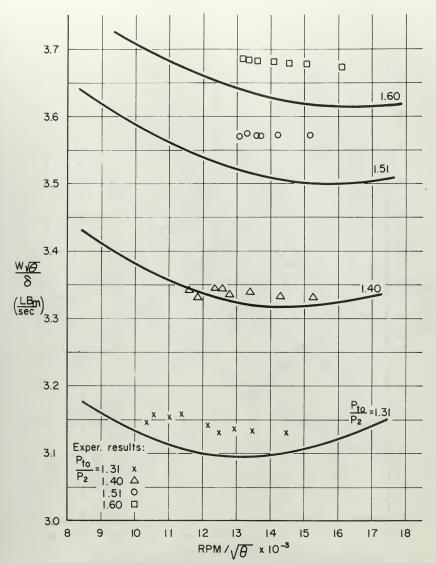
FIG. 24



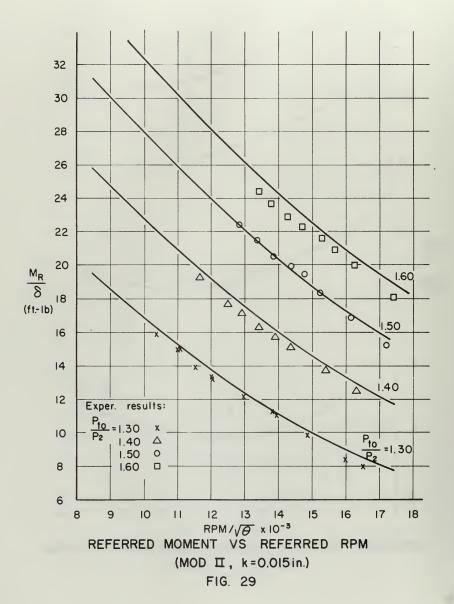


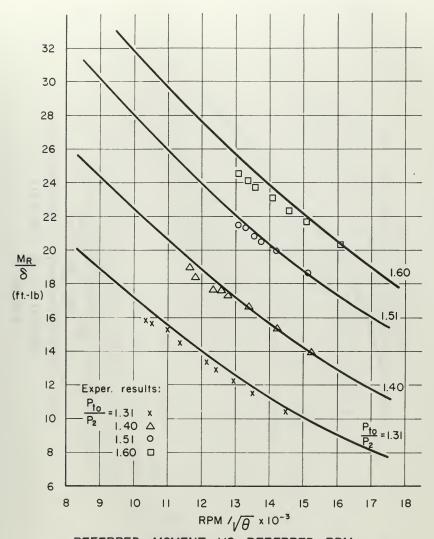


VARIATION OF REFERRED FLOWRATE WITH REFERRED RPM (MOD II, k=0.015 in.)
FIG. 27

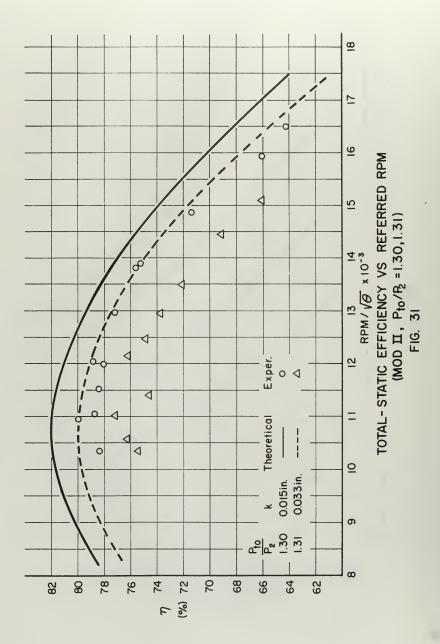


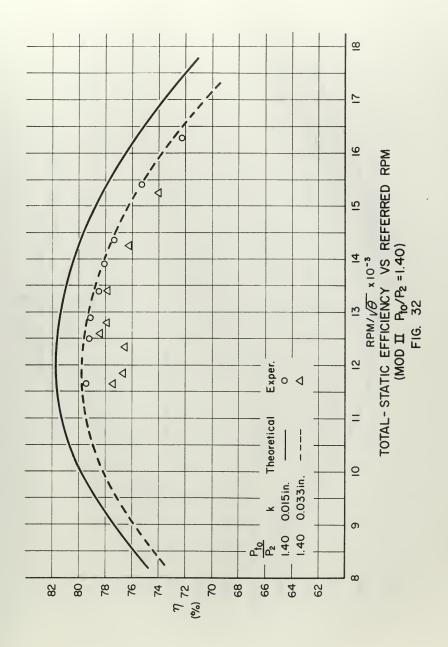
VARIATION OF REFERRED FLOWRATE WITH REFERRED RPM (MOD II, k=0.033 in.) FIG. 28

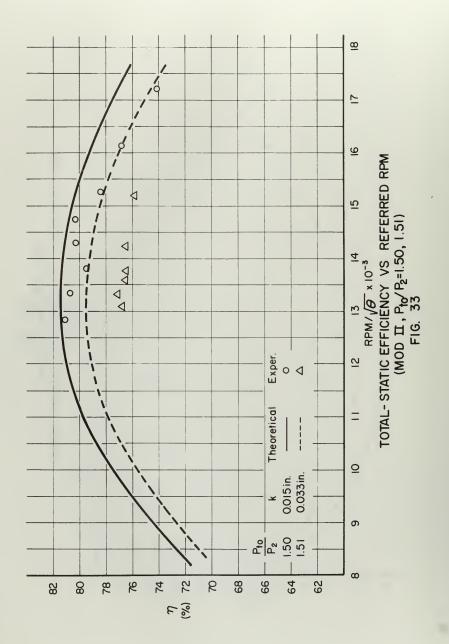


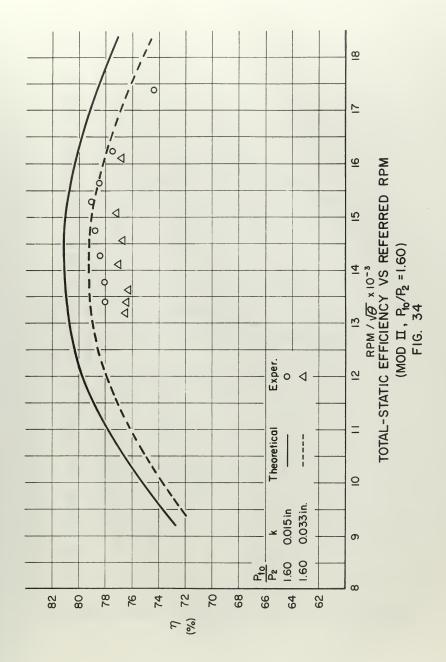


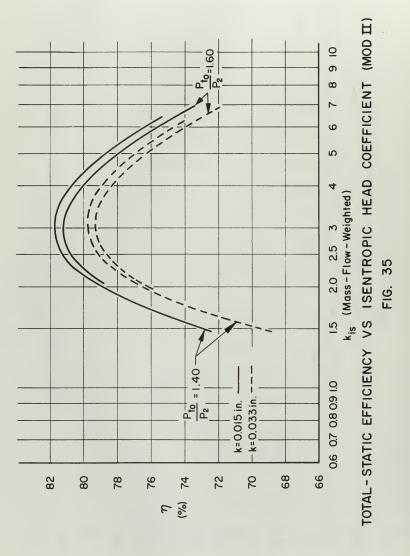
REFERRED MOMENT VS REFERRED RPM (MOD II, k=0.033in.)
FIG. 30

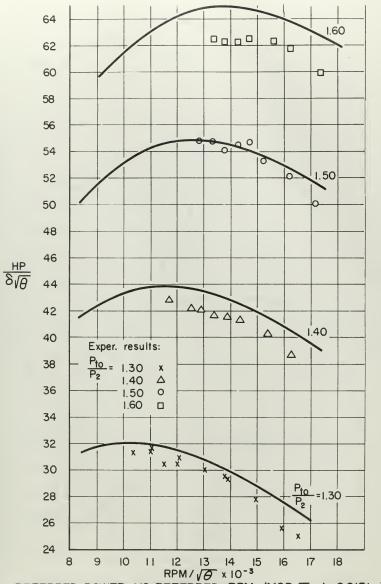




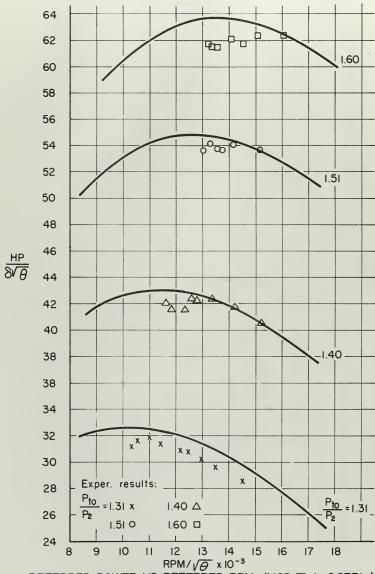




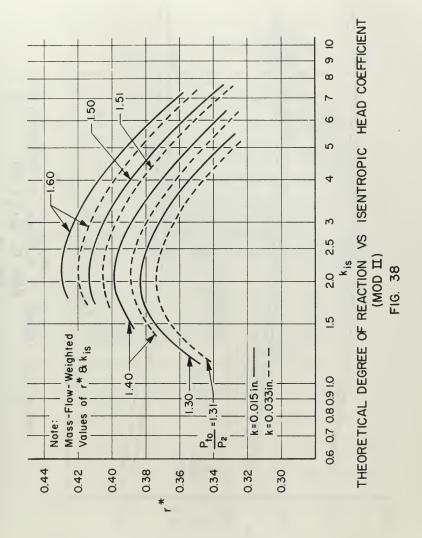


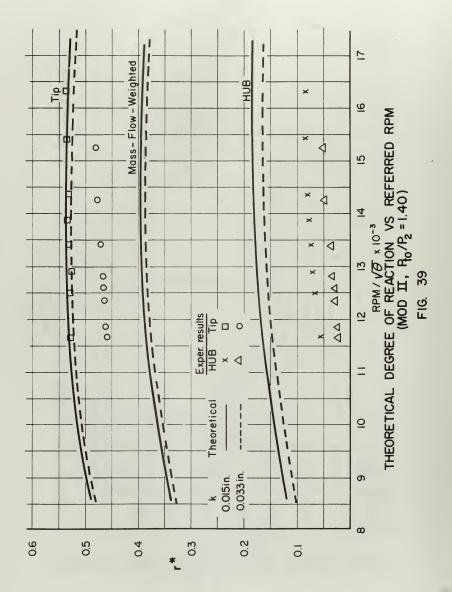


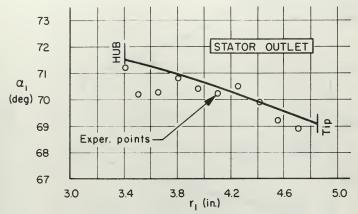
REFERRED POWER VS REFERRED RPM (MOD II, k=0.015in.)
FIG. 36



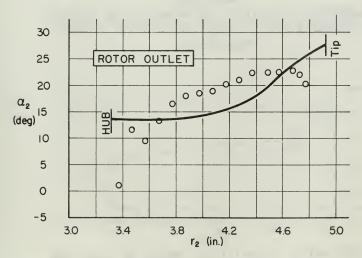
REFERRED POWER VS REFERRED RPM (MOD II, k=0.033in.) FIG. 37

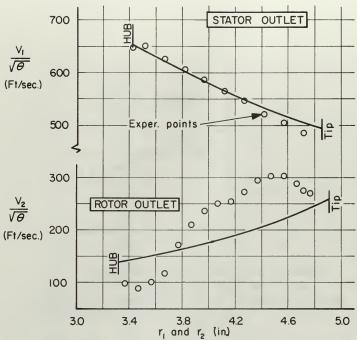




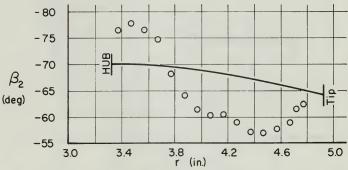


ABSOLUTE FLOW OUTLET ANGLES AS FUNCTION OF RADIUS (MOD II, k=0.015 in.,  $P_{10}/P_2$ =1.40, RPM/ $\sqrt{\Theta}$  = 13,934) FIG. 40

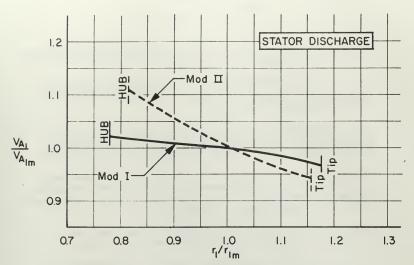




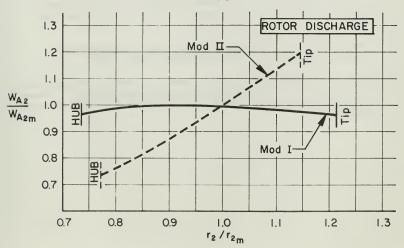
REFERRED VELOCITIES AS FUNCTION OF RADIUS (MOD II, k=0.015 in.  $P_{10}/P_2$ =1.40, RPM/ $\sqrt{\theta}$ =13,934)
FIG. 41



RELATIVE ROTOR FLOW OUTLET ANGLE AS FUNCTION OF RADIUS (MOD II, k=0.015in.  $R_{10}/P_2$ =1.40, RPM/ $\sqrt{\theta}$ =13,934) FIG. 42



PLOTS OF AXIAL VELOCITY RATIOS VS RADIUS RATIOS AT PEAK EFFICIENCY (MODI & II, Pio/P2 =1.40) FIG. 43



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#### APPENDIX A

### DEVELOPEMENT OF EQUATIONS

## 1. Equation of Motion 11

The equation of motion for relative flow is

$$\nabla H_{R} = \widetilde{W} \times (\nabla \times \widetilde{V} + 2\widetilde{\omega}) + T\nabla S$$
 (2)

In cylindrical coordinates:

$$\nabla H_{F} = \frac{1}{12} \frac{\partial H_{F}}{\partial \theta} + \frac{1}{12} \frac{\partial H_{R}}{\partial x} + \frac{1}{12} \frac{\partial H_{R}}{\partial y}$$
 (87)

$$\frac{1}{10} \left[ V_{1} \left( \frac{1}{10} \frac{1}{10} \right) - W_{1} \left[ \frac{1}{10} \frac$$

$$\overline{W} \times 2\overline{\omega} = \overline{i}_{c} (2\omega W_{c}) - \overline{i}_{g} (2\omega W_{c})$$
 (89)

$$T\nabla \varepsilon = T\left(\overline{i}_{\theta} + \overline{i}_{\theta} + \overline{i}_{\theta} + \overline{i}_{\theta} + \overline{i}_{r} + \overline{i}_{r} + \overline{i}_{r} + \overline{i}_{r} + \overline{i}_{r}\right)$$
(90)

Equating Eq. 87 to the components in Eqs. 88-90,

$$\overline{l}_{\theta} : \frac{1}{r} \frac{\partial H_{\theta}}{\partial \theta} = \frac{V_{10}}{r} \left[ \frac{\partial W_{10}}{\partial \theta} - \frac{\partial (rV_{10})}{\partial \theta} \right] - \frac{W_{10}}{r} \left[ \frac{\partial (rW_{10})}{\partial r} - \frac{\partial V_{10}}{\partial \theta} \right] - 2\omega W_{10} + \frac{1}{r} \frac{\partial s}{\partial \theta}$$
(91)

$$\frac{1}{z} \cdot \frac{\partial H_{\varphi}}{\partial r} = W_{r} \left[ \frac{\partial W_{r}}{\partial z} - \frac{\partial V_{r}}{\partial r} \right] - \frac{W_{r}}{r} \left[ \frac{\partial W_{h}}{\partial z} - \frac{\partial (rW_{r})}{\partial z} \right] + \frac{1}{r} \frac{s}{z}$$
(92)

$$\frac{1}{i}: \frac{\partial H_{R}}{\partial r} = \frac{W_{I}}{r} \frac{\partial (r W_{I})}{\partial r} - \frac{\partial W_{R}}{\partial \theta} - W_{R} \left[ \frac{\partial V_{IR}}{\partial z} - \frac{\partial W_{R}}{\partial r} \right] + 2 \omega W_{0} + T \frac{\partial s}{\partial r}$$
(93)

<sup>11</sup> Eckert, op. cit., pp. 149-155.

For assumed axisymmetric flow,  $\frac{\partial(}{\partial \partial} = 0$ ,  $\frac{\partial}{\partial \partial} = 0$ ,  $\frac{\partial}{\partial} = 0$ ,  $\frac{\partial}{\partial}$ (94)

$$\overline{i_z} : \frac{\partial H_R}{\partial z} = W_\Gamma \frac{\partial W_\Gamma}{\partial z} - W_\Gamma \frac{\partial W_R}{\partial r} + \frac{W_J}{r} \frac{\partial (rW_J)}{\partial z} + \frac{\partial s}{\partial z}$$
 (95)

$$\frac{1}{r} : \frac{\partial H_R}{\partial r} = \frac{W_D}{r} \frac{\partial (rW)}{\partial r} - W_A \frac{\partial W_r}{\partial z} + W_A \frac{\partial W_A}{\partial r} + 2\omega W_U + T \frac{\partial s}{\partial r}$$
(96)

$$\frac{\partial (rW)}{\partial z} = -\frac{W_r}{W_A} \frac{\partial (rW)}{\partial r} - 2\omega r \frac{W_r}{W_A}$$
(97)
Replacing  $\frac{\partial (rW)}{\partial r}$  in Eq. 95 by its equivalent in Eq. 97,

$$\frac{\partial H_R}{\partial H_R} = W_r \frac{\partial Z}{\partial W_r} - V_{V_r} \frac{\partial V_R}{\partial W_R} + \frac{VV_r}{V_r} \left( -\frac{W_r}{W_r} \frac{\partial (r V_r)}{\partial r} \right) - 2\omega \frac{W_r W_r}{W_R} + T \frac{\partial z}{\partial z}$$
(98)

$$\frac{\partial H_{R}}{\partial z} = -\frac{W_{2}V_{1}r}{rVv_{R}}\frac{\partial (rW_{2})}{\partial r} + W_{r}\frac{\partial W_{r}}{\partial z} - W_{r}\frac{\partial W_{R}}{\partial r} - 2\omega \frac{W_{2}r}{W_{R}} + T\frac{\partial z}{\partial z}$$
(99)

Multiplying Eq. 96 by  $W_r$  and Eq. 99 by  $W_a$ ,

$$W_{A} = \frac{1}{2} = \frac{W_{A} V_{A}}{C} = \frac{1}{2} \frac{\partial (V_{A})}{\partial V_{A}} + W_{A} W_{A} = \frac{\partial W_{A}}{\partial E} - W_{A} W_{A} = \frac{\partial W_{A}}{\partial V_{A}} + 2 \omega W_{A} W_{A} + W_{A} + \frac{\partial S}{\partial E}$$
(101)

Adding Eqs. 100 and 101,

$$V_V \frac{\partial H_R}{\partial V} + V_A \frac{\partial H_R}{\partial F} = \overline{I} \left( W_V \frac{\partial S}{\partial F} + W_A \frac{\partial S}{\partial F} \right)$$
 (102)

For adiabatic flow HD is constant along a streamline, therefore,

$$W_{at} \cdot \nabla H_{a} = 0 \tag{103}$$

For axisymmetric flow,

$$\overline{W}_{3}^{+} \cdot \nabla H_{F} = 0 = W_{4} \frac{\partial H_{R}}{\partial z} + W_{7} \frac{\partial H_{R}}{\partial r}$$
 (104)

From Eq. 104.

$$\frac{\partial H_R}{\partial x} = -\frac{W_C}{W_L} \frac{\partial H_R}{\partial x} \tag{105}$$

Substituting Eq. 105 into Eq. 102,

$$W_{r} = \frac{3r}{3H^{2}} + W_{r} \left( -\frac{\sqrt{N}}{M^{2}} \frac{3r}{3H^{2}} \right) = T \left( W_{r} \frac{3r}{3S} + W_{r} \frac{3z}{3S} \right)$$
 (106)

$$\frac{\partial S}{\partial z} = -\frac{W_f}{V_f} \frac{\partial S}{\partial z} \tag{107}$$

Using Eqs. 105 and 107, Eq. 99 becomes

$$-\frac{W_{r}}{W_{0}}\frac{\partial H_{R}}{\partial r} = -\frac{V_{c}V_{c}V_{c}}{cV_{c}V_{c}}\frac{\partial (rV_{c})}{\partial r} + W_{r}\frac{\partial W_{r}}{\partial z} - W_{r}\frac{\partial V_{c}}{\partial r} - 2\omega\frac{W_{c}V_{c}}{V_{c}A} - \frac{W_{r}}{W_{A}} + \frac{35}{3r}$$
(108)

Multiplying Eq. 108 by  $\left(-\frac{\sqrt{2}}{2}\right)$ , gives

$$\frac{\partial He}{\partial t} = \frac{VV}{V} \frac{\partial (VV)}{\partial r} - W_A \frac{\partial W_C}{\partial z} + VV_A \frac{\partial V}{\partial r} + 2 \omega W_A + 7 \frac{\partial S}{\partial r}$$
 (109)

Eq. 109 is identical with Eq. 96, therefore Eq. 109 is the equation that must be solved. With  $\left[VV_{\mu} \frac{\partial V_{\nu,k}}{\partial \gamma} = \frac{1}{2} \frac{\partial \left(V_{\nu,k}^{-2}\right)}{\partial \gamma}\right]_{2}$  Eq. 109 becomes

$$\frac{3(W_1)^2}{3r} = -\frac{3}{2} \frac{3(W_1)^2}{3r} + \frac{3(W_2)^2}{3r} + 4 \omega W_1 - 2 \frac{3He}{3r} + 2T \frac{3s}{3r} = 0$$
 (110)

Rewriting Eq. 110 for the rotor discharge and using the substitutions  $H_R = H_L = \frac{M_E^2}{2} \qquad \text{and} \qquad T_2 = \frac{H_E}{CP} = \frac{W_L^2}{2 CP} \text{ }$ 

$$\frac{3(V_{-}A)}{3V_{z}} - 2W_{A_{z}}\frac{3V_{C_{z}}}{3Z_{z}} + \frac{W_{z}}{C_{z}}\frac{3(C_{z}W_{0z})}{3C_{z}} + 4\omega W_{0z} - 2\frac{3}{3C_{z}}\left(H_{E} - \frac{U_{z}^{2}}{2}\right) + 2\left(\frac{H_{E}}{C_{P}} - \frac{W_{z}^{2}}{2C_{P}}\right)\frac{3c_{z}}{3C_{z}}$$
(111)

The last term of Eq. 111 can be expressed by

$$\frac{1}{C_{P}} \left[ -H_{E} - \left( V_{1} - V_{2} - V_{1} \right) \right] = \frac{1}{C_{P}} \left[ 2H_{E} - W_{0}^{2} - W_{A2}^{2} \left( 1 + \frac{W_{0}^{2}}{V_{A}^{2}} \right) \right] \qquad (112)$$
With  $\left( \frac{V_{C}^{2}}{V_{A}^{2}} - T_{1} - V_{2} \right)$  and  $\left( 1 - T_{A}N_{1} - N_{2} - \frac{1}{C_{2} + 2} \right)$  Eq. 112 becomes

$$\frac{1}{C_{P}} \left[ 2 H_{E} - V v_{z}^{2} \right] = \frac{1}{C_{P}} \left[ 2 H_{E} - W v_{z}^{2} \right] - \frac{W A_{z}^{2}}{C_{P} Cos^{2} \lambda_{z}}$$
(113)

Using Eq. 113, and after rearranging, Eq. 111 becomes

$$\frac{\partial (V_{11})}{\partial r_{2}} = V_{11} = \frac{\partial V_{11}}{\partial z} = \frac{W_{12}^{2}}{r_{1}} = \frac{\partial S_{2}}{\partial r_{2}} + 2 \frac{W_{02}}{r_{2}} = \frac{\partial (C W_{2})}{\partial r_{2}} + 4 \omega V_{1} = -\frac{\partial H_{E}}{\partial r_{2}} + \frac{\partial (U_{2}^{2})}{\partial r_{2}} + \frac{1}{C_{P}} \left[ 2H_{E} - W_{1}^{2} \right] \frac{\partial S_{2}}{\partial r_{2}} = 0$$
(114)

Noting that

$$\frac{\lambda(U_2^2)}{\lambda(U_2^2)} = \frac{\lambda(U_2^2)}{\lambda(U_2^2)} = 2\lambda L^2 r_2 \tag{115}$$

Eq. 114 is

$$-2\frac{3HE}{2T} + 2\pi \cdot C + \frac{1}{CH} \left[ 2HE - M_{\odot}^{2} \right] \frac{35z}{37z} = 0$$
(116)

This equation is made non-dimensional by multiplying by  $\left(\frac{\gamma_m}{\sqrt{r^2}}\right)$ where the subscript m refers to the mean streamline. sequent derivations the subscript 2 will be omitted.

$$\frac{r_{m}}{W_{A_{m}}^{2}} \frac{3(W_{h}^{2})}{3r} - 2\frac{W_{h}}{W_{A_{m}}^{2}} r_{m} \frac{3W_{h}}{3z} - \frac{W_{h}^{2}}{W_{h}^{2}} \frac{r_{m}}{c_{p} \cos^{2} \lambda_{2}} \frac{3s}{3r} + \frac{W_{0} r_{m}}{W_{A_{m}} r} \frac{3(\frac{r}{r_{m}} \frac{W_{0}}{W_{h}})}{3(\frac{r}{r_{m}})}$$

$$+4\frac{w_{n}^{2}W_{b}}{W_{h_{m}}^{2}}-2\frac{r_{m}}{v_{k_{m}}}\frac{\partial H_{c}}{\partial r}+\frac{2w^{2}r_{m}}{v_{k_{m}}^{2}}+\frac{r_{m}}{c_{p}}\left[\frac{2H_{c}}{W_{k_{m}}^{2}}-\frac{W_{b}^{2}}{W_{c}}\right]\frac{\partial S}{\partial r}=0$$
(117)

Now let 
$$\gamma = \frac{VV_{\hat{H}}}{VV_{L_{K_{n}}}}$$
 (118)

$$\gamma = \frac{\Gamma}{\Gamma_{m}}$$
 (119)

$$S^* = \frac{S}{Sp} \tag{120}$$

Equation 117 then becomes

$$+4\frac{U_{m}YTANB}{W_{F}} - \frac{2}{W_{h}^{2}} \frac{\partial H_{E}}{\partial A} + 2\frac{U_{m}U}{W_{h}^{2}} + \left[\frac{2H_{E}}{W_{h}^{2}} - Y^{2}TAN^{2}B\right]\frac{\partial S^{*}}{\partial X} = 0$$
 (121)

The fourth term in Eq. 121 can be expanded and rearranged as

$$=2Y^{2}TAN^{2}\frac{\partial TAN^{2}}{\partial x}+2YTAN^{2}\frac{\partial Y}{\partial x}+\frac{2Y^{2}}{x}TAN^{2}\beta$$
 (122)

with

$$\frac{\partial THNB}{\partial x} = \frac{1}{\cos^2 R} \frac{\partial B}{\partial x}$$

and

Equation 122 then is

$$\frac{\partial (Y^{2})}{\partial X} (1 + TAN^{2}\beta) - 2 \frac{Y^{2}}{W_{A}} Y_{M} \frac{\partial W_{f}}{\partial z} - \frac{Y^{2}}{\cos^{2}\lambda} \frac{\partial s^{*}}{\partial X} + 2 Y^{2} \frac{TAN\beta}{\cos^{2}\beta} \frac{\partial \beta}{\partial X} + \frac{2 Y^{2}}{X} TAN^{2}\beta$$

$$+ \frac{4 U_{M} Y TAN\beta}{W_{A_{m}}} - \frac{2}{W_{A_{m}^{2}}} \frac{\partial H_{E}}{\partial X} + 2 \frac{U_{M} U}{W_{A_{m}^{2}}} + \left[ \frac{2 H_{E}}{W_{A_{m}^{2}}} - Y^{2} TAN^{2}\beta \right] \frac{\partial s^{*}}{\partial X} = 0$$

$$\text{Multiplying by } \left( -\cos^{2}\beta / Y^{2} \right), \text{ and with } 1 + TAN^{2}\beta = \frac{1}{\cos^{2}\beta},$$

$$\frac{1}{Y^{2}} \frac{\partial (Y^{2})}{\partial X} + \cos^{2}\beta \left( -\frac{2 (m}{W_{A}} \frac{\partial W_{f}}{\partial z} - \frac{1}{\cos^{2}\lambda} \frac{\partial s^{*}}{\partial X} \right) + 2 TAN\beta \frac{\partial \beta}{\partial x} + \frac{2}{X} \sin^{2}\beta$$

$$+\frac{4U_{m} \sin \beta \cos \beta}{W_{Am} Y} + \frac{2U_{m}U\cos^{2}\beta}{W_{Am}^{2} Y^{2}} - \frac{2\cos^{2}\beta}{W_{Am}^{2} Y^{2}} \frac{\partial H_{E}}{\partial X} + \left[\frac{2H_{E}\cos^{2}\beta}{W_{Am}^{2} Y^{2}} - \sin^{2}\beta\right] \frac{\partial S}{\partial X} = 0$$
 (123)

The terms  $\cos^2 \lambda$  and  $\left(\frac{-2 \sqrt{M_A}}{M_A} \frac{\partial M_C}{\partial Z}\right)$  represent the effects due to streamline curvature and can be approximated by

$$-\frac{1}{N^{4}}\frac{9S}{9N^{4}} = +\frac{1}{N}\frac{\Gamma_{5}}{SL} \quad ; \quad K \approx 2$$
 (124)

$$\cos^2 \lambda = \frac{L^2}{L^2 + \left(\frac{\Delta R}{2}\right)^2} \tag{125}$$

If the streamline curvature in meridional planes is zero, the terms represented by Eqs. 124 and 125 will take on values of zero and one, respectively. With 5r positive as shown in Fig. 2, the plus sign is used for K at station 2 and the minus sign at station 1. The streamline slope as represented by Eq. 125 is the same at stations 1 and 2 since the streamline pattern is assumed to repeat itself after station 2.

Using Eqs. 124 and 125, Eq. 114 then becomes

$$\frac{d\left(\ln Y^{2}\right)}{dX} = -\cos^{2}\beta\left[\left(2Kr_{m}\frac{Sr}{L^{2}}\right)\left(\frac{L^{2}+\left(\frac{\Delta R}{2}\right)^{2}}{L^{2}}\right)\frac{dS^{*}}{dX}\right] - 2TAN\beta\frac{d\beta}{dX} - \frac{2}{X}SIN^{2}\beta$$

$$-\frac{4 \operatorname{Um} \operatorname{SIN} \beta \operatorname{cos} \beta}{\operatorname{Wam} Y} - \frac{2 \operatorname{Um} \operatorname{Ucos}^2 \beta}{\operatorname{Wa}_m^2 Y^2} + \frac{2 \operatorname{cos}^2 \beta}{\operatorname{Wa}_m^2 Y^2} \frac{d \operatorname{He}}{d \operatorname{W}} - \left[ \frac{2 \operatorname{He} \operatorname{cos}^2 \beta}{\operatorname{Wa}_m^2 Y^2} - \operatorname{SIN}^2 \beta \right] \frac{d s^*}{d \operatorname{W}}$$
(126)

For the enthalpy terms to be in terms of  $BTU/1b_m$ , they must be divided by 2gJ to keep the equation dimensionless. With  $C_1 = 2g$  J, Eq. 126 is

$$\frac{d(\ln Y^2)}{dX} = -\cos^2\!\!\mathcal{B} \left[ \left( 2\,\mathrm{KY_m} \frac{\mathrm{Sr}}{L^2} \right) - \left( \frac{L^2 + \left( \frac{\Delta R}{2} \right)^2}{L^2} \right) \frac{dS^*}{dX} \right] - 2\,\mathrm{TAN} \beta \, \frac{d\beta}{dX} - \frac{2}{X}\,\sin^2\!\!\beta \, \frac{1}{2} \left( \frac{1}{2}\,\mathrm{TAN} \beta \, \frac{d\beta}{dX} - \frac{1}{2}\,\mathrm{TAN} \beta \, \frac{d\beta}{dX} \right) \right] + \frac{1}{2}\,\mathrm{TAN} \beta \, \frac{d\beta}{dX} + \frac{1}{2$$

$$-\frac{4U_{m}\sin\beta\cos\beta}{W_{Am}Y} - \frac{2U_{m}U\cos^{2}\beta}{VV_{Am}^{2}Y^{2}} + \frac{C_{1}\cos^{2}\beta}{W_{Am}^{2}Y^{2}} \frac{dH_{E}}{dX} - \left[\frac{C_{1}H_{E}\cos^{2}\beta}{VV_{Am}^{2}Y^{2}} - \sin^{2}\beta\right] \frac{dS^{*}}{dX}$$
(127)

Noting the corresponding terms for the absolute flow of the stator, and with U=0, the equation for station 1 is

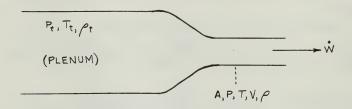
$$\frac{d(\ln Y^2)}{dX} = -\cos^2 \alpha \left[ \left( -K2r_m \frac{Sr}{L^2} \right) - \left( \frac{L^2 + \left( \frac{\Delta R}{Z} \right)^2}{L^2} \right) \frac{ds^*}{dX} \right] - 2 TAN \alpha \frac{d\alpha}{dX}$$

$$-\frac{2}{X} \sin^2 \alpha + \frac{C_1 \cos^2 \alpha}{Y^2 V_{Am}^2} \frac{dH}{dX} - \frac{C_1 H \cos^2 \alpha}{Y^2 V_{Am}^2} - \sin^2 \alpha \frac{ds^*}{dX} \tag{128}$$

### 2. Flow Function \$\overline{\phi}\$

Flowrate in  $lb_{m}/sec$  for the expansion process shown in Fig. 44 can be written

$$\dot{W} = \rho A V$$
 (129)



Expansion From Plenum

Fig. 44

Assuming an isentropic expansion,

$$\frac{1}{T_t} = \left(\frac{P}{P_t}\right)^{\frac{\chi}{\chi}} \tag{130}$$

and,

$$\overline{T}_{+} - \overline{T} = \overline{T}_{+} \left( 1 - \frac{T}{T_{+}} \right) = \overline{T}_{+} \left[ 1 - \left( \frac{F}{F_{-}} \right)^{\frac{N-1}{N}} \right]$$
 (131)

then using the thermodynamic relations

$$\frac{\sqrt{2}}{2gJc_F} = T_t - T \qquad \text{and} \qquad c_P = \frac{R}{J} \frac{\gamma}{\gamma - 1}$$
 (132)

the discharge velocity can be expressed by

$$V = \left\{ \text{agR} \frac{\delta}{\delta^{-1}} T_{t} \left[ 1 - \left( \frac{P}{P_{t}} \right)^{\frac{\delta^{-1}}{\delta}} \right] \right\}^{\frac{1}{2}}$$
 (133)

The density  $\rho$  is found with the isentropic assumption by

$$\frac{F_t}{P_t^{\ell}} = \frac{F}{P^{\ell}} \tag{134}$$

or

Using Eqs. 133 and 135, Eq. 129 becomes

$$\dot{W} = A \frac{P_t}{RT_t} \left( \frac{P}{P_t} \right)^{\frac{1}{\delta}} \left\{ 2gR \frac{8}{8-1} T_t \left[ 1 - \left( \frac{P}{P_t} \right)^{\frac{K-1}{\delta}} \right] \right\}^{\frac{1}{2}}$$
 (136)

This can be rearranged to give

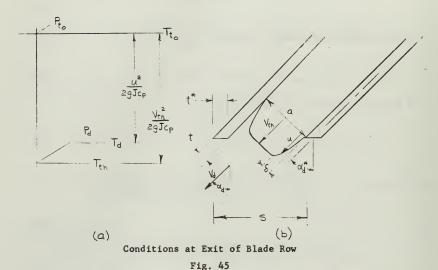
$$N = \frac{A P_{t}}{P_{t}} \left\{ \frac{28}{8-1} \left[ \left( \frac{P}{P_{t}} \right)^{\frac{2}{8}} - \left( \frac{P}{P_{t}} \right)^{\frac{8}{8}+1} \right] \right\}^{\frac{1}{2}}$$
(137)

or

$$\frac{\sqrt[4]{\sqrt{T_{\tau}}}}{4|P_{t}|}\sqrt{\frac{R}{r_{t}}} = \left\{\frac{28}{34}\left[\left(\frac{P}{P_{t}}\right)^{\frac{2}{8}} - \left(\frac{P}{P_{t}}\right)^{\frac{3+1}{8}}\right]\right\}^{\frac{1}{2}} \equiv \Phi$$
(138)

# 3. Restriction Factor & 12

The expansion process through a blade row is shown in Fig. 45 (a). The letter u represents the velocity at some point within the boundary layer and  $V_{\rm th}$  represents the velocity for an isentropic expansion. The various dimensions, angles, and velocities for the throat location are shown in Fig. 45(b).



The mass flowrate per unit blade height can be written

$$\dot{m} = \rho_{th} V_{th} \cos \alpha_d \left[ S - \frac{t}{\cos \alpha_d^2} - \frac{ZS}{\cos \alpha_d} \right] + \sum_{s} \int_{a}^{s} u \rho \, dy$$
where:

 $P_{th}$  = density at  $P_d$  for isentropic expansion (slugs/ft<sup>3</sup>)  $P_d$  = density corresponding to the velocity u (slugs/ft<sup>3</sup>)

 $\alpha_d^*$  = blade angle

δ = boundary layer thickness

<sup>12</sup> Vavra, M. H., Problems of Fluid Mechanics in Radial Turbo-machines (Rhode-Saint-Genese, Belgium: Von Karmán Institutéefor Fluid Dynamics, 1965) VKI Course Note 55b, pp. G46-50.

For Mach numbers of 0.8 or less, the discharge angle can be calculated quite accurately by

$$\alpha_{d} = \cos^{-1} \left[ \frac{a}{s - \frac{t}{\cos \alpha_{d}^{*}}} \right] \tag{140}$$

where:

a = minimum throat width

t = blade thickness

then with  $\eta = \frac{\gamma}{\zeta}$  and using Eq. 140 in Eq. 139,

$$\dot{m} = \rho_{th} V_{th} \frac{c\cos\alpha_d^*}{\cos\alpha_d^* - \tau} \left[ \frac{s\cos\alpha_d^* - t}{\cos\alpha_d^*} - 2s\left(\frac{s\cos\alpha_{d-1}^* t}{a\cos\alpha_d^*}\right) \right] + 2\rho_{th} V_{th} s \int_{th}^{t} \frac{u}{V_{th}} d\eta$$
(141)

which can be expressed as

$$\dot{m} = \rho_{\ell_h} V_{\ell_h} \alpha \left\{ 1 - \frac{8}{2} \left( 1 - \frac{1}{2} \frac{u}{V_{\ell_h}} d \eta \right) \right\}$$
 (142)

Assuming constant pressure at the throat,

and

$$\frac{P}{Reh} = \frac{T_{th}}{T} \tag{143}$$

where
$$\frac{T_{th}}{T} = \frac{T_{to} - \frac{V_{th}^2}{z_9 J c_p}}{T_{to} - \frac{u^2}{z_9 J c_p}} = \frac{T_{to} - \frac{V_{th}^2}{z_9 J c_p}}{T_{to} - \left(\frac{V_{th}^2}{z_9 J c_p}\right) \left(\frac{u}{V_{th}}\right)^2} = \frac{1 - \lambda_e}{1 - \lambda_e \left(\frac{u}{V_{th}}\right)^2}$$
In Eq. 144 the term  $X_e$  is
$$X_{-1} = \frac{1 - \lambda_e}{1 - \lambda_e} \left(\frac{u}{V_{th}}\right)^2 = \frac{1 - \lambda_e}{1 - \lambda_e}$$

$$X_{e} = \left| -\frac{T_{th}}{T_{to}} \right| = \left| -\left(\frac{P_{d}}{P_{to}}\right)^{\frac{\chi-1}{\gamma}} \right|$$
 (145)

The boundary layer is assumed to be turbulent and the profile can therefore be expressed by

$$\frac{u}{V_{th}} = \left(\frac{\gamma}{\delta}\right)^m = \gamma^m \tag{146}$$

Then with the displacement thickness of the boundary layer defined as

$$S^* \equiv S \left[ 1 - \left( 1 - \chi_e \right) \int_{1 - \chi_e \eta^{2m}}^{1 - \chi_e \eta^{2m}} d\gamma \right]$$
(147)

Eg. 142 can be written

$$\dot{m} = \rho_{th} V_{th} a \left[ 1 - \frac{\xi \delta^*}{a} \right] \tag{148}$$

The loss coefficient to the throat is

$$y' = 1 - \frac{V_d^2}{V_{th}^2}$$
 (149)

The loss coefficient can also be expressed in terms of the kinetic energy lost, by

$$\mathcal{G} = \frac{\Delta \dot{E}}{\dot{m} \frac{V_{ch}^2}{\Delta \dot{E}}} = \left| - \frac{\dot{E}}{\dot{m} \frac{V_{ch}^2}{\Delta \dot{E}}} \right|$$
 (150)

where E represents energy rate at the discharge due to the average velocity  $\mathbf{V}_{d}$ , or

$$\dot{E} = \rho_{th} V_{th} \left[ a - \xi \delta \right] \frac{V_{th}^2}{2} + \xi \int_0^{\delta} \left( u \rho \frac{u^2}{2} \right) dy \tag{151}$$

Equation 151 can be rewritten to give

$$\dot{E} = \int_{th} \frac{\sqrt{3}}{2} a \left\{ \left| - \left( \left| - \right| \times_{e} \right) \right| \frac{\eta^{3m}}{\left( \left| - \times_{e} \right|^{2m} \right)} d\eta \right\}$$
(152)

Then using the energy thickness which is defined as

$$S^{***} \equiv S \left[ 1 - \left( 1 - \chi_e \right) \int_{\left( 1 - \chi_e \right)^{2m}}^{1} d\eta \right]$$
 (153)

the loss coefficient can be expressed by

$$y = 1 - \frac{P_{th} \frac{V_{th}^{3}}{2} a \left[ 1 - \frac{S^{***}}{a} \right]}{\frac{1}{m} \frac{V_{th}^{2}}{2}} = 1 - \frac{P_{th} \frac{V_{th}^{3}}{2} a \left[ 1 - \frac{S^{***}}{a} \right]}{\frac{V_{th}^{3}}{2} a \left[ 1 - \frac{S^{**}}{a} \right]}$$

or

$$\mathcal{G} = 1 - \frac{1 - \mathcal{E} \frac{\xi^{****}}{\alpha}}{1 - \mathcal{E} \frac{\xi^{**}}{\alpha}} \tag{154}$$

The restriction factor g represents that part of the throat opening in which would occur the uniform theoretical velocity, therefore

$$\xi = 1 - 2 \frac{\xi^*}{a} \tag{155}$$

Now defining an energy parameter H\*\*\* by

$$H^{***} \equiv \frac{\S^{***}}{\S^*} \tag{156}$$

Eq. 155 becomes

$$\xi = \frac{H^{***} - 1}{H^{***} - 1 + \frac{1}{4}}$$
s coefficient 4 in Eq. 157 accounts for the losses that

The loss coefficient y in Eq. 157 accounts for the losses that occur from the inlet to the throat of the blade channel. No means exists for predicting this loss coefficient. Half the total loss coefficient for the blade row y will be used to represent these losses. Eq. 157 then becomes

$$\xi = \frac{H^{***} - 1}{H^{***} - 1 + \frac{1}{2}}$$
 (158)

4. Method of Evaluating H\*\*\* 13

By use of the binomial theorem, the denominator of the integral part of Eqs. 147 and 153 can be expanded yielding

The integral of Eq. 153 is then

Then, integrating and evaluating gives

$$\int_{1-X_{e}\eta^{2m}}^{1} d\eta = \frac{1}{3m+1} + \frac{X_{e}}{5m+1} + \frac{X_{e}^{2}}{7m+1} + \frac{X_{e}^{3}}{9m+1} + \frac{X_{e}^{4}}{11m+1} + \cdots$$

Equation 153 can now be expressed as

$$\frac{\S^{***}}{\S} = \left[ -\left( \frac{1}{3m+1} + \frac{\chi_e}{5m+1} + \frac{\chi_e^2}{7m+1} + \frac{\chi_e^3}{9m+1} + \frac{\chi_e^4}{11m+1} + \cdots \right) + \frac{\chi_e}{3m+1} + \frac{\chi_e^2}{5m+1} + \frac{\chi_e^3}{7m+1} + \cdots \right]$$

<sup>13</sup> Eckert, op. cit., pp. 159-160.

or,

$$\frac{5^{***}}{5} = 1 + (\chi_{e} - 1) \left( \frac{1}{3m+1} + \frac{\chi_{e}}{5m+1} + \frac{\chi_{e}^{2}}{7m+1} + \frac{\chi_{e}^{3}}{9m+1} + \frac{\chi_{e}^{4}}{11m+1} + \cdots \right)$$

also,

$$\frac{5^{***}}{5} = \left( \chi_{e^{-1}} \right) \left( \frac{1}{\chi_{e^{-1}}} + \frac{1}{3^{m+1}} + \frac{\chi_{e}}{5^{m+1}} + \frac{\chi_{e}^{2}}{7^{m+1}} + \frac{\chi_{e}^{3}}{9^{m+1}} + \frac{\chi_{e}^{4}}{11^{m+1}} + \cdots \right)$$

In a similar manner, the expression for Eq. 147 is

$$\frac{S^*}{S} = \left( \chi_{e-1} \right) \left( \frac{1}{\chi_{e-1}} + \frac{1}{m+1} + \frac{\chi_{e}}{3m+1} + \frac{\chi_{e}^2}{5m+1} + \frac{\chi_{e}^3}{7m+1} + \frac{\chi_{e}^4}{9m+1} + \cdots \right)$$

The energy parameter can now be expressed by

$$H^{***} = \frac{\frac{1}{X_{e-1}} + \frac{1}{3m+1} + \frac{X_{e}}{5m+1} + \frac{X_{e}^{2}}{7m+1} + \frac{X_{e}^{3}}{9m+1} + \frac{X_{e}^{4}}{11m+1}}{\frac{1}{X_{e}^{-1}} + \frac{1}{m+1} + \frac{X_{e}}{3m+1} + \frac{X_{e}^{2}}{5m+1} + \frac{X_{e}^{3}}{7m+1} + \frac{X_{e}^{4}}{9m+1}}$$
(159)

The number of terms in the numerator and denominator of Eq. 159 is considered sufficient to give good convergence for H\*\*\*.

Although the exponent m for the turbulent boundary layer is dependent on Reynolds number, it is taken to be a constant for this analysis.

### COMPUTATION OF OUTLET ANGLES AND LOSS COEFFICIENTS

### 1. Outlet Angles

The absolute discharge angles for the stator and the relative discharge angles for the rotor are computed by using the same methods. For these calculations, the assumption is made that the outlet angles are not influenced by the flow incidence angles. The factors which do affect the discharge angles are:

- Blade geometry; this results in the angles being a function of the radius since spacing and possibly profiles change from hub to tip.
- 2. Radial tip clearance; the effects of tip clearance are assumed to influence the rotor flow from the mean streamline out to the tip, with the largest effect being near the tip.
- 3. Exit Mach number; in accordance with the experimental results surveyed by Ainley, the outlet angles are a function of Mach number. Values of the angles are calculated for Mach numbers M of 0.5 and 1.0. Then a smooth curve of outlet angle versus M is drawn between these points with an inflection point at M=0.75. Below M=0.5, the flow angles are assumed to be equal to the value computed for M=0.5.

Vavra's formula is used for the first approximation of the stator outlet angles for M.≤0.5.

$$\alpha^* = \cos^{-1}\left(\frac{\alpha/s}{K_t}\right) \tag{160}$$

where

$$K_{t} = 1 - \frac{2.7}{10^{3}} \left(\frac{t}{5} \times 0^{3.3} \frac{a}{5}\right)$$
 (161)

The effects of blade curvature are next taken into account by using the method given by Ainley 14.

$$\alpha = \alpha^* + 4\left(\frac{s}{e}\right) \tag{162}$$

<sup>14</sup> Ainley, op. cit., pp. 3-4.

In Eq. 162, "e"is the mean radius of curvature of the upper surface of the blade profile between the passage throat and the trailing edge. This quantity is approximated by

$$e = \frac{j^2}{8b} \tag{163}$$

where j and b are shown in Fig. 46.

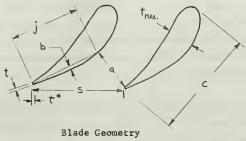


Fig. 46

The discharge angle for M=1.0 is found in a manner similar to that described by Ainley,

$$\alpha_{(M_1=1,0)} = \cos^{-1}\left(\frac{\alpha}{s-r^*}\right) \tag{164}$$

The change in  $ext{ } ext{ } ext{$ 

$$\delta \alpha_{T} = \delta \alpha_{H} = \alpha_{M_{(M_{i}=0.5)}} - \alpha_{M_{i}(M_{i}=1.0)}$$
(165)

Relative discharge angles for the hub and mean radii of the rotor are found by the same method as used for the stator. The effects of tip clearance are accounted for in the computation of  $\beta_2$  for the tip position. With Ainley's formula, the flow angle  $\beta_2$  for the tip is

$$\beta_{z} = TAN^{-1} \left\{ \left[ 1 - X \left( \frac{k}{h} \right) \left( \frac{\cos \beta_{o}}{\cos \beta_{z}'} \right) \right] TAN \beta_{z}' + X \left( \frac{k}{h} \right) \left( \frac{\cos \beta_{o}}{\cos \beta_{z}'} \right) TAN \beta_{o} \right\}$$
(166)

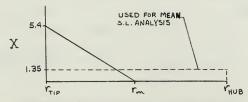
The new terms in Eq. 166 are:

h - blade height

 $\beta_2'$  - discharge angle before til clearance is accounted for

X - factor which depends on the type of shroud

Ainley suggests that X be set equal to 1.35 if no shroud is arranged at the rotor tips. However, Ainley's value has been determined with a mean streamline analysis. To localize the effects of tip clearance, X has been changed to 5.4, a value which gives about the same overall tip clearance effect as if the value of X=1.35 were used for the mean streamline. Fig. 47 shows that the areas under the two lines representing X as a function of radius are approximately equal. The slope of the solid line representing X for this method is obtained because of the assumed linear variation of outlet angle with radius between the mean radius and the tip.



Tip Clearance Factor X as Function of Radius

Fig. 47

Sample calculations for stator outlet angles are shown in Table I. A sample of rotor outlet angle computations can be seen in Table II.

TABLE I
SAMPLE CALCULATIONS FOR STATOR OUTLET ANGLES

# MOD I turbine

$M_1 \leq 0.5$			
radius - in.	3.597	4.125	4.950
spacing (s) - in.	1.738	1.994	2.392
throat opening (a) - in.	0.455	0.590	0.804
trailing edge thickness (t) - in.		- 0.044	
<u>a</u> <b>\$</b> .	0.2619	0.2960	0.3360
$\left(\frac{\mathbf{c}}{s_{3}}100\right)^{3.3} = A$	21.4	13.6	7.48
$K_t = 1 - (0.0027) \text{ (A) } (\frac{a}{s})$	0.9849	0.9891	0.9932
a sK <sub>t</sub>	0.2658	0.2991	0.3384
$\alpha^* = \cos^{-1}(\frac{a}{sK_t})$	74.6°	72.6	70.2
$4 \left(\frac{8}{e}\right)$	0.703	0.9106	1.37
$0 = 0 + 4(\frac{s}{e})$	75.3°	73.50	71.6

 $M_1 = 1.0$ 

Use 4.125 in. from centerline as representative radius.

Projected trailing edge thickness (t\*)-in. 0.13  

$$\alpha_{im} = \cos^{-1}(\frac{a}{s-t*})$$
 - deg. 71.5°  
 $\delta \alpha = 73.5 - 71.5 = 2.0$ 

TABLE II
SAMPLE CALCULATIONS FOR ROTOR OUTLET ANGLE

# SAMPLE CALCULATIONS FOR ROTOR OUTLET ANGLES MOD II turbine tip clearance = 0.033 in.

M ≤ 0.5

m - 0.5				
radius - in.	3.3235	4.197	4.918	
spacing (s) - in.	1.595	1.4643	1.7158	
throat opening (a) - in.	0.382	0.560	0.721	
trailing edge thickness (t) - in.		0.048		
<u>a</u> s	0.3295	0.3824	0.4202	
$\left(\frac{t}{s}\right)(00)$	4.1397	3.278	2.7975	
$\left(\frac{t}{s}   \infty\right)^{3.3} = A$	109.0	50.3	29.8	
$K_{t} = 1 - (0.0027)(A)\left(\frac{a}{5}\right)$	0.9030	0.9481	0.9662	
<u>a</u> 5 K+	0.3649	0.4033	0.4349	
$\beta^* = \cos^{-1}\left(\frac{a}{s K_1}\right)$	-68.5°	-66.2	-64.2	
4 ( <u>s</u> )	1.75	2.37	3.04	
$\beta = \beta^* - 4\left(\frac{s}{e}\right)$	<u>-70.3</u> °	<u>-68.6</u> °	-67.2°	
k - in.			0.033	
<b>k</b> /h			0.0207	
X=(4) (1.35)=type shroud factor			5.4	
cos Bo, Bo=30°			0.86603	
cos Bz', Bz' = -67.2°			0.38268	
$(X)(\frac{k}{h})\frac{\cos B_0}{\cos B_2} = B$			0.25297	
$ an eta_2'$			-2.4142	
$(1-B)\tan \beta_2$			-1.8035	
tan $\mathcal{B}_o$			0.57735	
$(1-B)\tan \beta_z' + B TAN\beta_o = C$			-1.6574	
$\beta_2 = \tan^{-1}(C)$			<b>-</b> 58.9	

# TABLE II (continued) SAMPLE CALCULATIONS FOR ROTOR ANGLE

Use radius = 4.197 in. as representative radius.

Projected trailing edge thickness (t\*)-in. 0.108 
$$\beta_{z_m} = \cos^{-1}\left(\frac{\alpha}{s-t^*}\right) -65.6^{\circ}$$
$$\beta = -68.6+65.6 = -3.0$$

radius - in. 3.3235 4.197 4.918  $\beta_2$  -67.3° -65.6° -62.6°

## 2. Loss Coefficients

Loss coefficients were predicted using Ainley's methods. The techniques used are completely described in Refs. 1 and 2. The assumptions, definitions, and equations which are necessary for a basic understanding of the method are described below.

It is assumed that loss coefficients are independent of Mach number and, for a given glade geometry, are a function of the flow incidence on the blade. Losses are divided into profile losses, secondary losses, and losses due to tip clearance. Mixing losses are not accounted for separately but are included in profile and secondary losses.

A parameter Y is defined by

For incompressible flow, which is assumed for these calculations, the loss coefficient is

$$\mathcal{L}_{j} = \frac{\mathcal{E}_{j}}{1 + \mathcal{E}_{j}}$$
The subscript  $j$  on Y in Eq. 168 indicates the type loss represented; i.e., Y would be the profile loss.

For zero incidence, the Y for profile losses is

$$Y_{P(i=o)} = \left\{ Y_{P(\beta_o=o)} + \left( \frac{\beta_o}{\beta_d'} \right)^2 \left[ Y_{P(\beta_o=-\beta_d')} - Y_{P(\beta_o=o)} \right] \right\} \left( \frac{t_{m/c}}{o.z.} \right)^{-\beta_o/\beta_d'}$$
(169)

where:

 ${{\cal B}_{\!\!d}}'$  - blade discharge angle not accounting for tip clearance  $t_m$  - maximum blade thickness

c - blade chord

Equation 169 is applicable to both the stator and rotor. However, for stator use the blade inlet angle  $\mathcal{S}_o$  is zero. The quantities  $Y_{P(\mathcal{S}_o = P_d')}$  and  $Y_{P(\mathcal{S}_o = C)}$  are taken from Fig. 4 in Ref. 2 where these quantities are plotted as a function of blade solidity. The maximum value of  $t_{m/c}$  is set equal to 0.25 for thick blades.

To determine profile losses for incidences other than zero, Fig. 6 in Ref. 2 shows  $\frac{\gamma_P}{\gamma_{P(i=o)}}$  versus  $\frac{i}{\sqrt{5}}$  where is defined as that incidence where the losses are twice the minimum losses.

Secondary and tip clearance losses are computed by

$$Y_{s} + Y_{K} = \left[\lambda + B \frac{k}{h}\right] \left[\frac{C_{L}}{s/c}\right]^{2} \frac{\cos^{2} B_{d}}{\cos^{3} B_{m}}$$
(170)

where:

 $\beta_{m}$  - the mean of the gas inlet and exit angles of the blade row  $\beta_{m}$  - factor representing type shroud, similar to X  $\frac{C_{L}}{s/c} = 2\left(\text{TAN}\beta_{o} - \text{TAN}\beta_{d}'\right)\cos\beta_{m} = \int (i)$   $\lambda_{m} = \int \left[\left(A_{d}/A_{i}\right)^{2}/\left(1 + \frac{\text{T.D.}}{\text{O.D.}}\right)\right]$ (171)

in Eq. 171:

 $A_d = (annulus area at discharge) (cos <math>\beta_d$ )

 $A_i = (annulus area at inlet) (cos <math>\beta_o$ )

I.D. and O.D. - inner and outer diameters, respectively, at the blade exit

Values of \( \) are obtained from Fig. 8 in Ref. 2. Ainley suggests a value of 0.5 for B for the type shroud used. With the same reasoning as previously used for the effect of tip clearance on discharge angles, B is set equal to 2.0 for the losses at the tip and zero for the mean radius and hub locations.

Sample calculations for the prediction of stator loss coefficients are shown in Table III. Sample calculations for prediction of stalling incidence and rotor loss coefficients are shown in Table IV.

### TABLE III

# SAMPLE CALCULATIONS FOR STATOR LOSS COEFFICIENT

MOD I turbine

radius = 3.597 in.

pitch (s) = 1.738 in.

inlet angle ( $\alpha_{\circ}$ ) = 0 outlet angle ( $\alpha_{\downarrow}$ ) = 75.3° chord (c) = 2.691 in.

All figures referred to are in Ref. 2.

$$Y_{P(i=o)} = \begin{cases} Y_{P(\alpha_i=o)} + \left(\frac{\alpha_o}{\alpha_i}\right)^2 \left[Y_{P(\alpha_o=-\alpha_i)} - Y_{P(\alpha_o=o)}\right] \left(\frac{t/c}{o.z}\right)^{-\alpha_o/\alpha_i} \\ \alpha_o = o & \frac{P_{ITCH}}{CHORD} = \frac{1.738}{2.691} = 0.6469 \end{cases}$$

$$Y_{s} = \left[ \Lambda + B \frac{k}{h} \right] \left[ \frac{C_{s}}{s/c} \right]^{2} \left[ \frac{\cos^{2} \alpha_{s}}{\cos^{3} \alpha_{m}} \right] \qquad k = 0$$

$$\lambda = \int \left[ \left( \frac{A_2}{A_1} \right)^2 / \left( 1 + \frac{1.D.}{0.D.} \right) \right] = \int \left\{ \left( \cos 75.3^{\circ} \right)^2 / \left( 1 + \frac{6.6}{9.898} \right) \right\}$$
 I.D. = 6.6 in.   
 0, D. = 9.898 in.

$$h = f(0.0386)$$
  $h = 0.0056$  (Fig. 8)

$$\alpha_m = TAN^{-1} \left( \frac{TAN 75.3^{\circ}}{2} \right) = TAN^{-1} (1.906) = 62.3^{\circ}$$

$$\frac{C_L}{s_{/C}} = 2 \left[ TAN \alpha_0 - TAN \alpha_1 \right] \cos \alpha_m = 2 \left[ -\left( TAN - 75.3^{\circ} \right) \right] \cos 62.5^{\circ}$$

$$\frac{C_L}{5/c} = 3.5438$$

$$Y_S = [0.0056][3.5438]^2 \left[ \frac{\cos^2 75.3^{\circ}}{\cos^8 62.3^{\circ}} \right] = 0.0451$$

$$Y_s + Y_p = 0.0929$$

CORRECTION FACTOR FOR TRAILING EDGE THICKNESS , F= 1.03 (Fig.9)

$$y = \frac{0.0949}{1+0.0949} = 0.0866$$

#### TABLE IV

# SAMPLE CALCULATIONS FOR STALLING INCIDENCE AND ROTOR LOSS COEFFICIENTS

MOD II turbine 
$$s = 1.7158$$
 in.  $\beta_{o} = 30^{\circ}$  radius = 4.918  $k = 0.033$  in.  $\beta_{z}' = -67.2^{\circ}$   $t_{max} = 0.728$  in.  $c = 1.967$  in. blade height(h) = 1.595 in. Figures referred to are in Ref. 2. 
$$Y_{P(i=o)} = \begin{cases}
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Y_{P(i=o)} + \begin{cases} \frac{\beta_{o}}{\beta_{z}'} \end{cases} & Y_{P(\beta_{o}=-\beta_{z}')} - Y_{P(\beta_{o}=o)} \\ \frac{\beta_{o}'}{\beta_{o}'} \end{cases} & Y_{P(\beta_{o}=-\beta_{o}')} & Y_{P(\beta_{o}=$$

TABLE IV (continued)
SAMPLE CALCULATIONS FOR ROTOR LOSS COEFFICIENT

Vis	-1.5	0.0	1.5
$i = (i/i_s) i_s$	-38.25	0	38.25
$\beta_i = i + \beta_0$	-8.25	30	68.25
TANB,	-0.1450	0.5774	2.5065
$\frac{TANB_1 + TANB_2' = A}{2}$	-1.2619	-0.9008	0.0638
Bm = TAN- (A)	-51.6	-42.0	3.6
cos Bm	0.6212	0.7430	0.9980
TANB, - TANBZ' = B	2.2339	2.9563	4.8854
C1/s/c=2(B) cos Bm	2.7752	4.3928	9.7510
[ 2/5/6] 2	7.7016	19.2964	95.0813
1/cos 1 Bm	4.1726	2.4385	1.0061
$Y_{S} + Y_{K} = (0.003906) \left[ \frac{C_{L}}{5/c} \right]^{2} \frac{1}{\cos^{8}\beta_{m}}$	0.2345	0.3434	0.6980
YP/YP(=0) (Fig.6)	2.1	1.0	4.5
YP = [YP/YP(1=0)] YP(1=0)	0.1390	0.0662	0.2979
YP + Ys + YK	0.3735	0.4096	0.9959
•	9) 0.3922	0.4301	1.0457
$y = \frac{C}{1+C}$	0.2817	0.3007	0.5112

#### APPENDIX C

### COMPUTER PROGRAM

A computer program using Fortran IV was written for use with the performance analysis. The program was used to predict performance values for the MOD I and MOD II turbines. However, an attempt was made to keep the program general. If the methods of this thesis are utilized for expressing turbine characteristics, this program can be used for other single stage axial turbines.

Input for the program consists of information representing the turbine characteristics, an indicator specifying the detail of output desired, and the conditions for which performance values are to be obtained.

Twenty input cards are used to introduce the turbine characteristics. The information contained on each of these cards, with required dimensions, is listed below.

- 1. Stator mean streamline gas outlet angles (ALFAM-radians) for mean streamline exit Mach numbers of 0.5, 0.7, 0.75, 0.8 and 1.0.
- Relative rotor mean streamline gas outlet angles (BETAMradians) for the same Mach numbers listed in (1.).
- 3. Radii ahead of the stator (RC-in.) for the five streamlines such that the flow area is divided into four equal parts.
- 4. Assumed radii of the five streamlines at the stator exit (RS-in.).
  - 5. Same as (4.), only for the rotor exit plane (RR-in.).
- The predicted stator loss coefficients (ZETAS) for the hub, mean radius, and tip.
- Ten values of incidence ratio (RINC) ranging from -2.0 to 1.6, in increments of 0.4.
- Ten values of rotor loss coefficients (ZETAR1) for the hub corresponding to the incidence ratios of (7.).
  - 9. Same as (8.), only for the mean radius (ZETAR3).
  - 10. Same as (8.), only for the tip (ZETAR5).

- 11. Length L shown in Fig. 2 (CL-in.) and the curvature factor K (CK).
- 12. Differences in stator gas outlet angles (DALF-radians) from the angle for the mean streamline, for streamlines 1, 3, and 5.
- Same as (12.), only for the relative rotor outlet angles (DBET).
- 14. Ten radii for the stator exit plane, equally spaced or approximately equally spaced, ranging from the radius of the hub to the radius of the tip (R1-in.).
- 15. Ten values of stator throat opening (Al-in.) corresponding to the radii of (14.).
  - 16. Same as (14.), only for the rotor exit plane (R2-in.).
- 17. Ten values of rotor throat opening (A2-in.) corresponding to the radii of (16.).
- 18. Inlet blade angles for the rotor (BETO-degrees) for the hub, mean radius, and tip.
- 19. Ten stall incidence angles (STALI-degrees) corresponding to the radii in (14.).
- 20. The radial tip clearance of the rotor (TIPC-in.), the number of stator blades (ZS), and the number of rotor blades (ZR).

Input card number 21 specifies the output to be printed, and its use will be described later. The remaining input specifies the conditions for which performance values are to be found, and enters the estimations of the flow Mach numbers used for the first approximations in the iteration process. Input card number 22 specifies the number of sets (NSETS) of operating conditions for which solutions are to be found. A card containing the following information is used for each point specified by NSETS:

- a. Estimated Mach number of the flow ahead of the stator (AMC).
- b. Estimated Mach number of the mean streamline flow after the stator (AMS).
  - c. Total inlet pressure (PTO-psi.).
  - d. Total inlet temperature (TTO-OR).

- e. Estimated relative Mach number of the mean streamline flow at the rotor exit (ANR).
  - f. The rotor speed (RFM).
  - g. The ratio of total inlet to static discharge pressure (PR).

There are eight subroutines in addition to the main or executive part of the program. The subroutines and their main functions are listed below:

- Subroutine PARAB is used to determine the coefficients of the parabolic equations used to approximate curves.
- Subroutine LSQPOL determines by the method of least squares the coefficients of the fourth order polynomials used to approximate curves.
- Subroutine CHAN computes the flowrate and the reference flowrate from conditions ahead of the stator.
- Subroutine STATOR determines the axial velocity ratios that satisfy the equation of motion at station 1.
- 5. Subroutine FLOWREF computes the reference flowrate at the stator exit and the rotor exit, and adjusts the axial velocity of the mean streamline flow at these locations to satisfy overall continuity.
- Subroutine SLINE checks streamline continuity for the stator exit and the rotor exit, and determines new streamline radii to satisfy streamline continuity.
- Subroutine ROTOR1 converts the absolute flow properties ahead of the rotor to relative flow properties.
- 8. Subroutine ROTOR2 determines the axial velocity ratios at station 2 that satisfy the equation of motion.

The main program computes the mass-flow-weighted value of total inlet to static discharge pressure ratio, and adjusts the Mach number of the flow ahead of the stator to obtain the pressure ratio specified. The overall turbine performance values of  $\gamma$ ,  $M_R$ , H.P.,  $r^*$ , and  $k_{ie}$  are also computed in the main program.

Computer output representing the solution for each set of conditions is printed on two pages. The flow properties at the stator exit plane are printed on one page. A sample of this output can

be seen in Table V. The second page contains the flow properties for the rotor as well as the overall turbine performance values. A sample of the second page of output is in Table VI. Most of the output quantities are self-explanatory. However, the following symbols may not be readily recognized and are defined as:

SLINE - streamline

ZETAP - one-half the total loss coefficient, used for calculating  $\boldsymbol{\xi}$ 

A - throat opening of the blade channel

W-FRAC - fraction of the flowrate w between the hub and the indicated streamline

BO - rotor blade inlet angle

INCID - rotor flow incidence angle

PSIR - relative flow velocity coefficient  $\frac{W_2}{W_{2th}}$ 

In addition to the main output there are additional WRITE statements in the program. The additional output allows the user to follow more closely the intermediate steps of the solution process. This output is also helpful when program changes are made which require debugging. The extra output is not desirable for production runs because of the large amounts of computer running time and printout that result. When the additional output is wanted, 1 is read into the computer on input card number 21 for the indicator IND. If 0 is read into the computer for IND, only the main output representing the final solution will be printed.

The statements of the program are listed on the following pages.

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12,714 223 223	C, TTO, RPM, B 9.2, F8.2, F NE R1 11), X1(11), Z1 3F8.5, F8.4	11 * SQRT( 17 * 3 157 * 3 17 * 17 * 1	VA1 (I), VU1(	SLINE R2 2)/(49.01*SQRT(T2 8.3,2F8.2,2F8.5, I.RR(I),X2(I),T
175) 712, 712 1750, 123, 22 1750, 123, 22 1750, 223, 22 1750, 223, 22 171)	SLI RS(	SHIZIWH C	SLINE SLINE 1 1 VAI 9 1 1 F8	-IANHIL
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	P 2	(1)	3, F8.2, F8.	VU2	(1), v2(1)		**************************************	111	7 7 7 7 7 7 7		ER2(L)-WF			PT1
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IYR(I), VA2(I)	WRITE (6,23) FORMAT (766H 9)	DO 190 I=1,5 ALFA2(I)=ATAN V2(I)=VA2(I)/ ALFA22(I)=ALF	BEIAZZ(I) = BEI FORMAT(I4, F9. WRITE(6,24)I.	FORMAT (768H STALI INCIE	FORMAT(14, F8. WRITE(6, 26)1,	DO 230 I=1,5 DELH(I)=(U(I) TT2(I)=TTQ-DE	PT2(I)=P2(I)* PTI(I)=P1(I)* T2IS(I)=TT0*(I)	ETAI(I)=(TT0-	ETAR(I)=(TTE( RSTAR(I)=(TTE) AKIS(I)=(TII	PSIR(I)=SQRT( DELH(IO)=0. DO 240 I=1,4	DELH(10)=DELH(10)+.5*(WP) HP=DELH(10)*CJ*WLBM/550. AWGH=HP*S50.(DMFG THFTA=SQRT(TTO/518.4)	DELTA=PTO/14. HP1=HP/(THETA AMOMI=AMOM/DEL	WLBM1=WLBM*THI	FORMAT (751H DO 250 I=1,5 FORMAT(14,6FR
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VAR 12=VAR 1**2

VAR 22=VAR 3**2

VAR 32=VAR 3**VAR 22-VAR 3**VAR 21**(VAR 22-VAR 32)*VAR 12**(VAR 32**VI)

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SUBROUTINE CHAN (TTO. AMC. PTO. RC, WLB W, WCHAN, WPERO)

COMMON IND RC(10), WPERG(10)

COMMON IND RC(10), WPERG(10)

VC=470/(1.+.2\*AMC\*\*2)

VC=49.01\*SQRT(TC)\*AMC

PC=FTG/(1.+.2\*AMC\*\*2)

RHG=PC/(53.35\*TC)\*\*2-RC(1)\*\*2)

AREA = 3.14f(\*\*C(5)\*\*2-RC(1)\*\*2)

WCHAN-WLBM/(PTG\*SQRT(32.174/(53.35\*TTO)))

WPERG(1) = 0.25

WPERG(2) = 0.25

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*353*353
5*((aLFAI(N)-ALFAI(I))/(X(N)-X(I))+(ALFAI(I)-ALFAI(*))/
))
C9=C9+1.

36 D 352 1=1.

DD 354 1=1.

DD 154 1=1.

DD 154
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SUBROUTINE FLOWR (PRAT, ZETAP, X, MI, PTE, PTO, TTE, TTO, AS, ZS, RS, AR, ZR, REAL& BAN, WCHAN, VA, WPER, CODE, WLBM, B, R, TIPC, A)

OD IMPERION PRAT(10), SETAP(10), R(10), M(10), PTE(10), TTE(10),

COMMON WINDER(10), B(20), A(10), R(10), M(10), M(10), PTE(10), TTE(10),

COMMON WINDER(10), B(20), A(10), R(10), R(10), M(10), M(10), PTE(10), TTE(10),

COMMON WINDER(10), B(20), A(10), R(10), R
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/TTO) *ARAT*XI*P4I = . F6.41
/2.
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1416*R(5)*TIPC/(ZR*AR*RR*(X(5)-X(4)))
16,416
XI,PHI,ARAT
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                         1+B[2]*R[I]+B[3]*R[I]**2+B[4]*R[I]**3+E

15*412*415

15*414*415

14.0**1416*R[5]*TIPC/(ZR*AR*RR*(X[5]-X

14.0**1416*R[5]*TIPC/(ZR*AR*RR*(X[5]-X

14.0**17*PII-ARAI

15**17*PII-ARAI

16**17*PII-ARAI

16**17*PII-ARAI

16**17*PII-ARAI

16**17*PII-ARAI
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VA(3)
PUTED-REQUIRED,AX VEI,2FIO.4,FIC.2
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X(4),WPER(4)
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2),X(3),WPER(3),
)) *(X(5)-X(4))/2.
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428
S*RS)
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1+SUM2)/WSUM
1+SUM2+SUM3)/WSUM
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C-WSUM)
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720,73C
                              ~ N
                             SAA
       700
           702
                   708
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DSDX1)
FC2=FC1**2

RRF=R(1)+B(2)*RRF+B(3)*RRF**2+B(4)*RRF**3+B(5)*RRF**4

IF(IND-1)732,721,721

21 IF(IND-1)732,721,729

24 WRITE (6,724)

25 FORMAT (747H SLINE XNEW HENEW DHEDX

26 FORMAT (14,F5,4,F9,2,F9,4,F9,6,F9,6)

27 FORMAT (14,F5,4,F9,2,F9,4,F9,6)

28 WRITE (6,726) I,X(II)*HE(II)*DHEDX(II)*S(II)*DSDXI(II)

30 CODE=40.

32 RETURN

END
                                                                721
729
724
                                                                                                                          726
                                                                                                                                                                       730
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SSUBROUTINE RCTOR!(VUI, VAI, RPM, U, BETAI, HE, TTE, PTE, XZ, PI, TI, WI, WUI, XX, BS, ZETAR, ZETAPR, RR, DHEDX, DSDX, S, UZ, OMEG, BPI, BR2, BR3, FSI, FS2, BEL&R B3, B4, B5, B6, RSI, RS3, RS 5, BETO, STALI, BINC)  REAL ** B3, B4, B5, B6  DDIMENSION VUI(10), VUI(10), VUI(10), BETAI(10), HE(10), TTE(10), PTE(10), TTE(10), TTE(	CCMMON IND C=2.*32.174*778.16*.24 OM G52 I = 1.74*778.16*.24 U [ ] = OM G6*RS [ 1 / 1.2 U [ ] = U [ ] = V [ 1 / 1 / 1 / 1 ] WU [ [ ] = V [ ]   V	HE(I)=TTE(I)*24-77-77-77-77-75-75-75-75-75-75-75-75-75-	RING=-2.0 GO TO 506 If (RING-16) 510, 504,504 WING=1.6 WING=1.6 FORMAT (/58H RATIO OF INCIDENCE TO STALL INCIDENCE FXCEEDED, USE IN 1/15 - F5.1) + B3(1) + B3(2) *RINC+83(3) *RINC+82(4) *RINC**3+B3(5) *RINC	1**4 2FTA(3)=B4(1)+B4(2)*RINC+B4(3)*RINC**2+B4(4)*RINC**3+B4(5)*RINC* 1*4 1*4 2FTA(5)=B5(1)+B5(2)*RINC+B5(3)*RINC**2+B5(4)*RINC**3+B5(5)*RINC 1**4 1**4 1**4	GETAR(1)=ZETA(1)+(RS(1)-RS1)/(RS3-RS1)*(ZETA(3)-ZETA(1))  GETAR(1)=ZETA(3)  GO TO 518  ZETAR(1)=ZETA(3)+(RS(1)-RS3)/(RS5-RS3)*(ZETA(5)-ZETA(3))  ZETAR(1)=ZETAR(1)/2.0  CONTINUE
		501	00000000000000000000000000000000000000		514 514 518 520
00002	00000000000000000000000000000000000000	0000	00022 00022 00025 00025 00026	0030	00000 00000 000000 000000 000000 000000

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DSDX(1)=(S(2)-S(1))((XZ(2)-XZ(1))

DSDX(2)=0.5*(DSDX(1)+(S(3)-S(2))/(XZ(3)-XZ(2)))

DSDX(3)=0.5*(DSDX(1)+(S(3)-S(2))/(XZ(3)-XZ(2)))

DSDX(4)=0.5*(DSDX(1)+(NZ(4)-XZ(3)-XZ(3)))

DHEDX(1)=(HE(2)-HE(1))/(XZ(2)-XZ(2))/(XZ(3)-XZ(2)))

DHEDX(2)=0.5*(DHEDX(1)+(HE(3)-HE(2))/(XZ(3)-XZ(2)))

DHEDX(3)=0.5*(HE(3)-HE(2))/(XZ(3)-XZ(2))+(HE(4)-HE(3))/

DHEDX(5)=(HE(5)-HE(4))/(XZ(5)-XZ(4))

DHEDX(5)=(HE(5)-HE(4))/(XZ(5)-XZ(4))

CCNTINUE

RETURN
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12. VUZ. VZ.
12. PZ. PRATZ
-, CK. DR. R.
              10) DSDX2(10) VA2(10)
VR(10) ZETAR(10)
10) SR2(10) VOLO(10)
PR T2(10) TZS(10)
10) DELR(10) RIS(10)
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AA, SR, TTE, PTE
SB2, RRF, DELR,
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(I)-RR3)/(RR1-RR3)*DBET(I)
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(I), I=1,5)
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-KYCLO(1)/YR(1))**2

-SK2(1)

)*(YGLO(1)/YR(1))**2

)*(YGLO(1)/YR(1))**2

(1)/YR(1))**2
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I)+DSDX2(I))*(SIN(I)
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+(X2(3)-X2(1)
+(X2(3)-X2(2)
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NM4

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m4500000

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BETA2(N)-BETA2(I))/(X2(N)-X2(I))+(BETA2(I)-BETA2
M))
                           7,F10.5,5F8.4)
) I,INDSI,DSDX2(I),R12(I),RI3(I),R14(I),RI(I),YR(I)
                                                                               3) *DBET(1
                                                                                         ETA(3)+(R(I)-RR3)/(RR5-RR3)*DBET(5
                                                                                              2)-x2(1))
5)-x2(4))
              ()-YR(I))
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                                                                                                               AN(BETA2(I))
          140,119
1120
1180
1180
1180
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DBETDX(I)

I(M))/(X2(

DO 225 I=

TAN I=-2.*
            224
       110
                                          70
                                                                             221
                                                                                    222
                                                                                         223
                                              251
0128
0128
0129
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ETA2(I))/(VA2(3)**R(I))
2/(VA2(3)**2*YR(I)**2)
                                                                                                                                                                                                                                                                                      )))**2/(C*HE(I))
3)*YR(I))**2*D4EDX(I)
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                 -ZETAR(1))
AM/(GAM-1.))
                  42 (1) *SIN(BE
18 TA2(1)) *SIN(BE
                                                                                                                                                                                                                                                                                            BETA2(1)
:2/(VA2(3
            /X2(I)
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                              \*
\*
SINI=-2.*SIN(BETAZ(I))**?/XZ(SRI(I)=-4.*U(3)*U(1)*COS(BETAZ(I))
SR2(I)=-4.*U(3)*U(1)*COS(BETAZ(I))
SR2(I)=-4.*U(3)*U(I)*COS(BETAZ(I))
SR2(I)=-4.*U(3)*U(I)*COS(BETAZ(I))
AA(I)=-4.*U(3)*U(I)*COS(BETAZ(I))
AA(I)=-4.*U(3)*U(I)*COS(BETAZ(I))
AA(I)=-4.*U(3)*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1)=-4.*U(1
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                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                   149
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			PTO/PI 1.25339 1.1579 1.1365	
TABLE V	0. 1.30	0.3801 0.5082 0.6161 0.7123	M-FRAC 0.0 0.2507 0.4998 1.0000	00000 00000 00000 00000 00000 00000 0000
	00 1200	1.01316 1.00784 1.00000 0.98457 0.96925	345.58 397.84 441.94 481.73 518.31	29941 2010-3 152-1 166-6 169-3
TATCR S	20 568.	ZETAPS 0.04520 0.032620 0.03810 0.03810	BETA1 60.35 42.91 16.90 -33.20	260.1 136.9 136.9 -28.2 -92.7
	NG. PTC	ZETAS 0.09040 0.08411 0.07924 0.07620 0.07340	AL FA1 76.27 74.61 73.28 71.60	623 5553 5623 5604 775 8
	SET	X1 0.782 0.900 1.000 1.090	13 72 14.42 14.42 15.13	6005 5005 6005 6005 6005 6005 6005 6005
		R1 4,229 4,5600 4,969	545.45 546.54 546.54 559.16 559.16	VA1 148.1 143.9 143.9
		SLINE 1. 22 33 54	SLINE 2 3 4	SLINE 22 24 54

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			VA2 129.7 138.9 142.8 149.2	0.3515 0.3562 0.3731 0.3639	1NC1D -20.009 -9.009 -9.009 -9.009 -9.009			LEMISEC	
	н	0	V 90.9085 0.9731 1.0000 1.0054	W-FRAC 0.25023 0.250211 0.75143 1.00000	STALI 13.48 34.80 40.42 38.93		0.07775 0.05974 0.03721 0.01286	EF RPM 1464.1	RSTAR 0.40041
	TABLE VI	M PR	ZETAPR 0.068447 0.06060 0.05456 0.08354	312.01 375.17 428.18 476.66 516.22	80 62.44 42.12 20.84 -1.61 -24.96	53.0.15 529.35 528.35 528.65 528.65	AKIS 4-1729 3-1729 2-551 1-825	ENT R	AKI S 2.689
1	SOLUTION	0 1200	ZETAR 0.16894 0.12121 0.10912 0.16709 0.21749	BETA2 -67.104 -69.578 -71.654 -71.104	######################################	5227 5227 5226 5226 527 196 196 196 196 196 196 196 196 196 196	RSTAR 0.13338 0.31729 0.43791 0.50116	E REF MOM	O NOI
	ROTCR SO	0 568.	114-129 114-239 15-25-3	ALFA2 2.2C9 0.9711 12.304 20.116	- 4444-1-1-1-1-1-1-1-1-1-1-1-1-1-1-1-1-1	20000000000000000000000000000000000000	PSIR 0.9374 0.9439 0.9126 0.9126	POWER 32.0	M-STAT
-		1 NO. PTO	541.27 5547.35 5530.55 553.46	PTO/P2 1.300 1.300 1.304 1.2999 1.2999	1229.8 1138.9 152.7	228833 28422 2842 2	0.8310 0.8388 0.8909 0.8329 0.7825	LORATE R BM/SEC 3-342	PT0/P2 1.3003
		SET	0.329 0.8729 1.0000 1.113	13.23 13.23 13.23 13.25 13.25 13.25 13.25	× × × × × × × × × × × × × × × × × × ×	PTI- 16.84 17.000 17.003	ETAS 0.99159 0.9238 0.9238	ENT FE	ETA 8412
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3. ABSTRACT						
The method of turbine perform	nance prediction	develo	ped by Vavra and			
Eckert has been refined in this ar						
of the three-dimensional calculate	-					
clearance effects on blade outlet angles and loss coefficients have been						
localized rather than averaged over		ght. A	n approximation			
for streamline curvature has been	used.					

Performance curves were determined for two single stage axial-flow turbines located at the Propulsion Laboratory of the Naval Postgraduate School. Test results were available for one of the turbines. Agreement between predicted and experimental performance values was generally within 3 per cent.

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